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Dielectric Constant and Density of Aqueous Alkali Halide Solutions by Molecular Dynamics: A Force Field Assessment

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The concentration dependence of the dielectric constant and the density of 11 aqueous alkali halide solutions (LiCl, NaCl, KCl, RbCl, CsCl, LiI, NaI, KI, CsI, KF, CsF) is investigated by molecular simulation. Predictions using eight non-polarizable ion force fields combined with the TIP4P/ ε water model are compared to experimental data. The influence of the water model and the temperature on the results for the NaCl brine is also addressed. The TIP4P/ ε water model specifically improves the accuracy of dielectric constant predictions compared to the SPC/E water model. The solution density is predicted well by most ion models. Almost all ion force fields qualitatively capture the decline of the dielectric constant with increasing concentration for all solutions and with increasing temperature for NaCl brine. However, the sampled dielectric constant is mostly in poor quantitative agreement with experimental data. These results are related to the microscopic solution structure, ion pairing and ultimately the force field parameters. Ion force fields with excessive contact ion pairing and precipitation below the experimental solubility limit generally yield higher dielectric constant values. An adequate reproduction of the experimental solubility limit should therefore be a prerequisite for further investigations of the dielectric constant of aqueous electrolyte solutions by molecular simulation.

Keywords: molecular simulation; dielectric constant; density; electrolyte solutions; force fields; ion pairing

I. INTRODUCTION

The unusually high dielectric constant of water under standard conditions of temperature and pressure ($\varepsilon = 78.5$)¹ is one of its anomalies,² which partly explains the dissolution of 1:1 salts in liquid water at low and moderate concentrations into single ions by preventing their aggregation into crystals.³ A substantial drop of the dielectric constant, usually caused by increasing temperature or salt concentration, may lead to the formation of ion pairs as a stable species in aqueous alkali halide solutions.⁴

Despite its significance, the dielectric constant of even the simplest aqueous monovalent salt solutions is poorly studied by experiment and sufficiently available only for 11 of the 20 alkali halide brines.^{5–19} In fact, the dielectric constant is not directly measurable in the laboratory.²⁰ It has to be extrapolated from the frequency-dependent dielectric constant in a non-trivial procedure, which requires elaborate measurement instruments, dielectric relaxation models and a complex estimation of model parameters.^{21–23}

An alternative to access this property, which was widely employed for water in the last three decades, is molecular modeling and simulation. The quality of the underlying force field models for solvent and solutes plays the central role for the accuracy of predictions with this approach. However, an extensive computational effort is required for a statistically sound sampling of the dielectric constant even for pure water. $^{24-26}$

Water and ions are primarily modeled by nonpolarizable Lennard-Jones (LJ) sites with superimposed point charges. The temperature dependence of the dielectric constant of water²⁷ or the recently found anomalously low dielectric constant of confined water²⁸ was qualitatively reproduced with these simple models. However, the quantitative accuracy of such predictions is rather disappointing.

Commonly employed water potentials in studies of a queous electrolyte solutions, such as SPC/E ($\varepsilon = 70.7$ at 298 K)²⁹ or TIP4P ($\varepsilon = 53$ at 293 K)³⁰, underestimate the dielectric constant. The TIP4P/2005³¹ water model, despite its superiority over other non-polarizable potentials,³² misses the dielectric constant by about 25 % ($\varepsilon = 58$ at 298 K). The general failure of non-polarizable water models in dielectric constant simulations was essentially attributed to the application of the same partial charges to empirically fit two distinct water molecule properties (potential energy and dipole moment).³³

In the search for optimal force field parameters for alkali halide ions, low concentration experimental data have mostly been used in fitting procedures.^{34–41} For example, the hydration free energy of the ions at infinite dilution was a popular choice among many groups.^{34–39} In an attempt to achieve reasonable predictions for the concentration range where cation-anion interactions become important, crystal properties^{37,42–44} or activity derivatives^{39,42,43} were considered. For a more compre-

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hensive discussion of common parametrization strategies employed for molecular modeling of aqueous electrolyte solutions, the interested reader is referred to the review of Nezbeda *et al.*⁴⁵

In general, however, these strategies yield unsatisfactory predictions for solution properties, such as density⁴⁶, water self-diffusion coefficient⁴⁷ or salt activity^{48,49}. In fact, the underprediction of the solubility limit of alkali halide salts,^{45,50–54} resulting in excessive ion pairing^{55–57} and/or premature crystallization,^{58–60} is a well-known problem of non-polarizable ion force fields. The cause may lie in unbalanced force field parameters.^{58,61} As crystallization is a slow process that takes place on time scales much longer than those typically accessed with molecular simulation, it is possible that many reported simulation results actually do not correspond to equilibrium states.

An accurate prediction of basic solution properties, such as density, does not only serve as an indication for an adequate representation of the microscopic structure,⁶² but is also necessary to correctly predict other properties.⁶³ The lack of experimental dielectric constant data, the inaccuracy of water models for dielectric constant predictions and generally inadequate ion force fields may explain why the dielectric constant of monovalent salt solutions was rarely studied with molecular simulation. After all, the prospect of dielectric constant investigations has recently improved dramatically by the development of water force fields parametrized to the dielectric constant (such as $TIP4P/\varepsilon^{64}$ or $TIP4Q^{65}$), methods for implicitly accounting for polarizability of non-polarizable models (such as charge scaling 66) and sufficient computational power for long simulations that are required for the dielectric constant sampling. A more detailed review of current trends in molecular modeling and simulation of aqueous electrolyte solutions was provided by Smith *et al.*⁶⁷

There are only a few studies concerning dielectric constant calculations of aqueous alkali halide solutions by simulation.^{43,68–72} Ion models used in these studies generally predict the decrease of the dielectric constant with increasing salt concentration qualitatively well, but underestimate the experimental values quantitatively. For example, Pethes⁷⁰ simulated the dielectric constant, density and self-diffusion coefficient of water and ions using 29 ion force fields for LiCl up to the experimental solubility limit of the according brine. In her study, most of the ion force fields overestimated the experimental dielectric constant at moderate concentrations. In a very recent study of aqueous NaCl solutions, a reduced dielectric constant decline at higher concentrations for some ion models was linked to excessive contact ion pairing.⁷²

The few comprehensive studies in the literature on molecular simulation of the dielectric constant of aqueous electrolyte solutions cover only one or a few salts. No study of the dielectric constant covering a wide range of alkali halide salts in aqueous solutions, together with a comparison of several different ion force fields, has been presented to date. The main objective of this paper is thus to provide such a comprehensive comparison. The performance of eight non-polarizable ion force fields combined with the TIP4P/ ε water model was assessed with respect to the prediction of the dielectric constant and density of 11 aqueous alkali halide solutions (LiCl, NaCl, KCl, RbCl, CsCl, LiI, NaI, KI, CsI, KF, CsF). These salts were chosen as they are the only ones out of the 20 alkali halides for which reliable experimental data are available for comparison.

The present approach is as follows: first, the role of the underlying water model was investigated by comparing dielectric constant predictions for aqueous NaCl solutions at 298 K, using different ion force fields together with the TIP4P/ ε or the SPC/E water model. As TIP4P/ ε yielded better predictions for all investigated ion force fields, all subsequent studies were carried out with that water model. Second, the temperature dependence of the dielectric constant was addressed, again using aqueous NaCl solutions as an example. Third, a comprehensive assessment of eight ion force fields for the alkali halides was carried out. All 11 salts for which experimental data are available for comparison were considered. Finally, the predictions for the dielectric constant were linked to ion pairing and salt solubility in light of the ion force field parameters.

II. METHODOLOGY

A. Molecular force fields

Throughout this work, ions and water molecules were represented by pairwise additive, non-polarizable force fields of the LJ 12-6 type with superimposed point charges. The interaction energy between particles i and j is then given by

$$u_{ij} = 4\varepsilon_{ij} \left[\left(\frac{\sigma_{ij}}{r_{ij}} \right)^{12} - \left(\frac{\sigma_{ij}}{r_{ij}} \right)^6 \right] + \sum_{l=1}^{N_i^\varepsilon} \sum_{m=1}^{N_j^\varepsilon} \frac{1}{4\pi\varepsilon_0} \frac{q_l q_m}{r_{lm}},$$
(1)

where σ_{ij} and ε_{ij} are the LJ parameters for size and energy, respectively, r_{ij} the site-site distance, q_l and q_m magnitudes of point charges of the ions or water molecules separated by the distance r_{lm} , N_i^e the total number of charges of particle *i* and ε_0 the vacuum permittivity. All interactions between unlike LJ sites were specified by the Lorentz-Berthelot (LB) combining rules

$$\sigma_{ij} = \frac{\sigma_{ii} + \sigma_{jj}}{2},\tag{2}$$

$$\varepsilon_{ij} = \sqrt{\varepsilon_{ii}\varepsilon_{jj}}.\tag{3}$$

In most of the present simulations, water-water interactions were described with the TIP4P/ ε water model.⁶⁴ This force field is a rigid body with a single LJ interaction site for the oxygen atom, two positive charges for the hydrogen atoms and one negative charge placed on a virtual site M on the H-O-H bisector. The force field parameters of the TIP4P/ ε and the SPC/E water models are given in Table S.I (Supplementary material). The LJ size and energy parameters of the TIP4P/ ε differ only slightly from those of the TIP4P/2005 force field. $TIP4P/\varepsilon$ was parametrized to reproduce the experimental density maximum of liquid water and the dielectric constant at ambient temperature. The significantly lower charge magnitude $q_{\rm H} = 0.527$ e and virtual site offset $d_{\rm OM} = 0.105$ Å of this model compared to TIP4P/2005 (0.5564 e, 0.1546 Å) result in an excellent coverage of the dielectric constant over a wide range of temperature and pressure. Both models give good predictions for most liquid state and vapor-liquid equilibrium properties of water.

In the present work, ions were modeled by one LJ site and a single point charge of fixed magnitude +1 e (cations) or -1 e (anions) in their center. Various ion force fields were studied, which differ by a factor of up to two for the size parameter σ and by up to four orders of magnitude for the energy parameter ε , as shown in Tables S.II and S.III. A summary of parametrization properties and water models employed in the fitting procedures for the ion models is given in Table I. It should be noted that in the original works of Jensen and Jorgensen³⁵, Reif and Hünenberger³⁶ and Gee *et al.*^{42,43} geometric combination rules were employed for the unlike LJ parameters.

As none of the ion force fields studied in the present work was adjusted explicitly for the use with the $TIP4P/\varepsilon$ water model, some remarks are given in the following about the transferability of these ion models to other water models than the ones employed in the fit. Transferability of the RDVH model parameters, fitted to reproduce the reduced liquid solution density, to four other water models was shown in the original study.⁴⁰ Moreover, these model parameters were used successfully for density predictions of methanolic⁷³, ethanolic⁷⁴ and water+methanol⁷⁵ alkali halide solutions. Consistent ion force fields from Gee et al. showed similar results for various properties with different water models. Mao and Pappu developed solvent-independent force field parameters that were fitted to crystal lattice properties.⁴⁴ Ion models from Joung and Cheatham parametrized with the TIP4P-Ew⁷⁶ water model showed a good compatibility with the TIP4P/2005 force field for aqueous $NaCl^{77}$ and LiCl⁷⁸ solutions due to the similar LJ geometry and parameters of TIP4P-family models. Hence, combining different ion model sets with the TIP4P/ ε water model appears to be reasonable.

B. Simulation details

An extended version of the molecular simulation tool $ms2^{79-81}$ was employed for the present calculations. Simulations of dielectric constant and density of electrolyte solutions were performed in the isothermalisobaric (NpT) ensemble by molecular dynamics. Eleven alkali halide brines (LiCl, NaCl, KCl, RbCl, CsCl, LiI, NaI, KI, CsI, KF, CsF) were sampled at 298 K, 1 bar and three concentrations for which experimental data are available, using eight different ion model sets together with TIP4P/ ε water. The remaining nine alkali halide salts (LiF, LiBr, NaF, NaBr, KBr, RbF, RbBr, RbI, CsBr) were not considered due to the lack of experimental data for comparison.

A cubic simulation volume with periodic boundary conditions contained a total of 1000 ion and solvent particles throughout. The composition of the solution is given in terms of the true ion mole fraction

$$x_{\rm ion} = \frac{n_{\rm C}}{2n_{\rm C} + n_{\rm W}},\tag{4}$$

where $n_{\rm C}$ and $n_{\rm W}$ are the mole numbers of cations and water, respectively. For the alkali halide brines of this study, the true ion mole fraction ranges from 0.007 to 0.095 mol mol⁻¹.

The pressure was kept constant by the Andersen barostat with a piston mass of $2.2 \cdot 10^9$ kg m⁻⁴. Velocity scaling was used to control the temperature. Following an equilibration of $5 \cdot 10^5$ time steps, production runs of 10^7 time steps were carried out. Simulation runs over more than 5-6 ns are needed for convergence of the dielectric constant.²⁴ Newton's equations of motion were solved with a Gear predictor-corrector scheme of fifth order with a timestep of $\Delta t = 1.2$ fs. The average runtime of a single simulation was eight hours on 4 compute nodes with 40 parallel cores per node on two Intel Xeon Gold 6148 CPU.

Electrostatic long-range contributions were calculated by Ewald summation with a real space convergence parameter $\kappa = 5.6$, up to 10 k-vectors in each Cartesian direction and conducting boundary conditions. The LJ cut off distance and the real-space cutoff were equal with 13.8 Å.

Statistical uncertainties were estimated by means of the block averaging method by Flyvbjerg and Petersen⁸² with a block length of 5000 time steps. Visual snapshots of molecular configurations were recorded every 5000 time steps.

Moreover, the radial distribution function (RDF) of several salts in water (NaCl, KCl, KI and LiI) was sampled to investigate the microscopic structure of the solutions. The number of contact ion pairs (CIP) for a 1:1 electrolyte solution can be computed from the cationanion RDF $g_{+,-}(r)$ with

$$n^{\rm CIP} = 4\pi \rho x_{\rm ion} \int_0^{r_{\rm min}} g_{+,-}(r) r^2 dr, \qquad (5)$$

where $\rho = N/V$ is the particle number density of the studied solution and r_{\min} the position of the first minimum of the cation-anion RDF.

In order to characterize the deviations between simulation results and experimental data for a given brine, the mean relative error (MRE) was calculated by

Ion models	Abbr.	Parametrization properties	Water	Combining
			model	rule
Horinek et al. ³⁴	$\operatorname{HMN-S^a}$	Hydration free energy, hydration entropy	SPC/E	-
Fyta <i>et al.</i> ^{38,39}	$Fyta^{b}$	Hydration free energy, activity derivatives	SPC/E	LB
Jensen-Jorgensen ³⁵	JJ	Hydration free energy, ion-water RDF	TIP4P	geometric
Reif-Hünenberger ³⁶	$RH-L^{c}$	Hydration free energy	SPC, SPC/E	geometric
Joung-Cheatham ³⁷	$\rm JC-T^d$	Hydration free energy, lattice energy, lattice	SPC/E, TIP3P,	LB
		constant, binding energy, ion-water RDF	TIP4P-Ew	
Gee <i>et al.</i> ^{42,43}	KBFF	Activity derivatives, ionic radii, lattice	SPC/E	geometric
		constant, ion-water RDF		
Reiser $et \ al.^{40,41}$	RDVH	Relative density, ion self-diffusion	SPC/E	LB
		coefficient, ion-water RDF	·	
$Mao-Pappu^{44}$	MP	Lattice energy, ion-ion distance	-	LB

TABLE I. Summary of ion force fields considered in this work with water models used for their parametrization.

^a Three parameter sets with different LJ energy parameters were given in the original work. The one with the smallest energy parameter was studied here. This ion force field does not provide parameters for the Rb⁺ cation.

^b This ion force field does not provide parameters for the Li⁺ and Rb⁺ cations. A modified LB combination rule was employed for

cation-anion interactions in the original work. ^c Three parameter sets based on different experimental hydration free energies of the proton and the SPC/E water model were given in the original work (L_E , M_E , H_E). The one for the lowest experimental hydration free energy (L_E) was used here.

the original work (L_E, M_E, H_E) . The one for the lowest experimental hydration nee energy (L_E)

^d TIP4P-Ew water model specific ion parameters were used in the present work.

$$MRE = \frac{1}{n} \sum_{i=1}^{n} \frac{|z_{\rm sim} - z_{\rm exp}|}{z_{\rm exp}},$$
(6)

where n is the number of data points per ion force field.

C. Sampling of the dielectric constant

The dielectric constant of non-polarizable fluids was computed from fluctuations⁸³ of the total dipole moment $\mathbf{M} = \sum_{i=1}^{N} \boldsymbol{\mu}_i$ in the simulation volume by

$$\varepsilon = \frac{4\pi}{3k_B T \langle V \rangle} \left(\left\langle \mathbf{M}^2 \right\rangle - \left\langle \mathbf{M} \right\rangle^2 \right) + 1, \tag{7}$$

where k_B is the Boltzmann constant and μ_i the individual dipole moment vector of the molecule. The angled brackets denote the ensemble average.

Long simulations bring the average dipole moment $\langle \mathbf{M} \rangle$ to zero,²⁵ as expected in the absence of an external field,⁸⁴ so that the term $\langle \mathbf{M} \rangle^2$ in Eq. (7) vanishes.

III. RESULTS AND DISCUSSION

In order to evaluate the performance of the studied ion and water models with respect to dielectric constant and molar density predictions, simulations were performed as described in Section II for a total of 327 state points. The results of 261 simulations were compared to experimental dielectric constant and density data in Figures 1 to 5. For the remaining simulations, indications of supersaturation or precipitation were found, as discussed further below. Tabulated numerical values of simulated properties together with their statistical uncertainties are listed in Tables S.IV to S.IX (Supplementary material).

The experimental data show that with rising salt concentration, the dielectric constant of the solution first decreases linearly (for dilute solutions) and then nonlinearly (for higher concentrations).⁸⁵ This phenomenon is called dielectric decrement and is primarily explained by the reduction of orientational correlation between water molecules in the proximity of ions.^{69,86} It should be pointed out that some experimental data were either delivered by a single research group or may be in disagreement with other available data. Therefore, some caution regarding the reliability of these data should be taken.

The specific density of all aqueous alkali halide solutions is an increasing linear function of ion mole fraction at a given temperature.⁸⁷ For many brines, the molar density shows a similar trend. Exceptions to this are the alkali iodides and CsCl, for which the high molar mass of at least one of the ions causes the molar density to decrease with concentration.

In order to investigate the sensitivity of the ion force fields to a change of the water model, aqueous NaCl solutions were studied, combining each of the ion force fields with the SPC/E or with the TIP4P/ ε water model. Results for the dielectric constant and the molar density under ambient conditions of temperature and pressure are shown in Figure 1 for three salt concentrations.

The experimental dielectric constant of pure water ($\varepsilon = 78.5$)¹ was predicted excellently with TIP4P/ ε ($\varepsilon = 77.4 \pm 1.5$), but underestimated with SPC/E ($\varepsilon = 69.0 \pm 1.3$). Both results are in good agreement with previous studies.^{3,29,64}

Regarding the concentration dependence, all ion models follow the same qualitative trends for both the di-

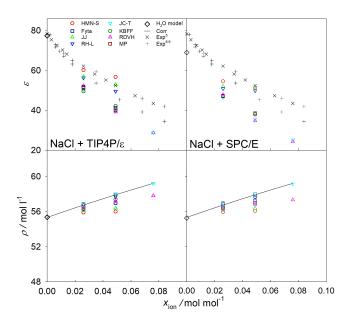


FIG. 1. Dielectric constant (top) and molar density (bottom) over true ion mole fraction of NaCl at 298 K and 1 bar. Simulation results for eight ion force fields combined with the TIP4P/ ε (left) and SPC/E (right) water models are represented as open colored symbols as shown in the legend. Experimental data for comparison are represented by cross symbols⁵. Other experimental data^{6–9} are depicted by plus symbols. The prediction for pure water is given by a diamond symbol. A correlation⁸⁷ to experimental density data is shown as a black line. Statistical uncertainties are within symbol size.

electric constant and the molar density, independent of the employed water model. The role of the water model becomes more significant when considering the quantitative agreement with experimental data. Simulations with TIP4P/ ε water led to a higher dielectric constant by around 8 % for all ion models at the lowest concentration. This is not surprising, since water-water interactions influence the dielectric constant to a greater extent in dilute solutions and the dielectric constant of pure TIP4P/ ε water is higher than that of pure SPC/E water.

The HMN-S, RH-L and JJ ion models generally led to a higher dielectric constant than other ion force fields. For more concentrated NaCl brines, the remaining ion models yielded an increase of up to 10 % for the dielectric constant with the TIP4P/ ε model and were thus in better agreement with the experimental data. The molar density predictions with all ion models were quite similar for both water models.

The use of the TIP4P/ ε model resulted in higher dielectric constant predictions for all ion force fields. These results are in line with previous findings, which suggest that 1-propanol⁸⁹ and NaCl⁸⁸ force fields combined with TIP4P/ ε water yielded dielectric constant predictions that were closer to experimental data. Thus, in the following, only TIP4P/ ε was considered as water model.

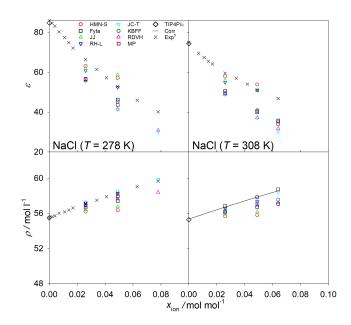


FIG. 2. Dielectric constant (top) and molar density (bottom) over true ion mole fraction of NaCl at 278 K (left) and 308 K (right) and 1 bar. Simulation results for eight ion force fields combined with the TIP4P/ ε water model are represented as open colored symbols as shown in the legend. Experimental data for NaCl brine⁵ are represented by cross symbols. The prediction for pure water is given by a diamond symbol. A correlation⁸⁷ to experimental density data is shown as a black line. Statistical uncertainties are within symbol size.

In order to check the experimentally observed decrease of the dielectric constant with rising temperature, NaCl brines were simulated for two additional temperatures (278 K and 308 K). The simulation results are compared with experimental data⁵ and a correlation⁸⁷ in Figure 2. For the lower temperature, experimental density data from Ref. 5 are shown instead of a correlation.

As in Figure 1 for aqueous NaCl solutions at 298 K, the same qualitative trends were observed at 278 K and 308 K for all ion force fields. An increase of the temperature reduced the discrepancy between simulations and experiments for the HMN-S, RH-L and JJ force fields, but raised deviations for the remaining ion models.

The same qualitative trends were observed for the performance of the ion force fields, irrespective of the temperature or the employed water model. Thus, in the following comprehensive study of alkali halides, only the TIP4P/ ε water model was used and only the temperature of 298 K was considered. The choice of this temperature was further motivated by the fact that for many of the investigated brines, experimental data are available only at that temperature.

The results of a comprehensive study of eleven alkali halides (LiCl, NaCl, KCl, RbCl, CsCl, LiI, NaI, KI, CsI, KF, CsF) in aqueous solution are presented in Figures 3 to 5. The following discussion summarizes the findings obtained by comparing the performance of the ion force

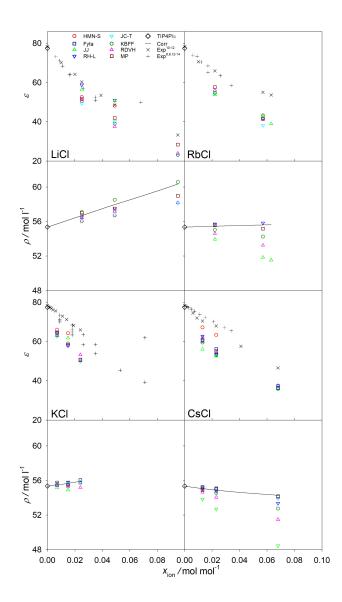


FIG. 3. Dielectric constant (top) and molar density (bottom) over true ion mole fraction of alkali chloride brines at 298 K and 1 bar. Simulation results for eight ion force fields combined with the TIP4P/ ε water model are represented as open colored symbols as shown in the legend. Experimental data for comparison for LiCl¹⁰, RbCl¹¹, KCl¹² and CsCl¹¹ brines are represented by cross symbols. Other experimental data for LiCl^{6,9,13}, RbCl^{9,13}, KCl^{6,9,13,14} and CsCl¹² brines are depicted by plus symbols. The prediction for pure water is given by a diamond symbol. A correlation⁸⁷ to experimental density data is shown as a black line. Statistical uncertainties are within symbol size.

fields in predicting the dielectric constant and molar density of these solutions.

For nearly all considered salts, the ion force fields are able to reproduce the experimental dielectric constant and density trends at least qualitatively. However, strikingly different results regarding the slope of the dielectric constant decrement were observed for various force fields.

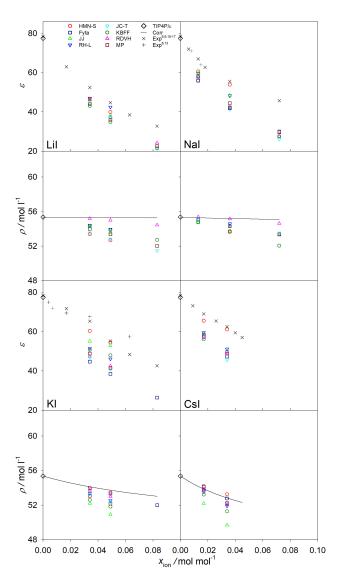


FIG. 4. Dielectric constant (top) and molar density (bottom) over true ion mole fraction of alkali iodide brines at 298 K and 1 bar. Simulation results for eight ion force fields combined with the TIP4P/ ε water model are represented as open colored symbols as shown in the legend. Experimental data for comparison for LiI¹⁵, NaI^{6,9}, KI¹⁶ and CsI¹⁷ brines are represented by cross symbols. Other experimental data for NaI¹³ and KI^{6,13} brines are depicted by plus symbols. The prediction for pure water is given by a diamond symbol. A correlation⁸⁷ to experimental density data is shown as a black line. Statistical uncertainties are within symbol size.

In most cases, the dielectric constant was underpredicted in comparison to the experimental data.

The molar density, with deviations below 1.6 % for almost all ion models, is in good agreement with experimental data. The HMN-S model set excellently reproduced the molar density of all potassium and cesium halide brines (except for CsF), but yielded the overall worst predictions for the remaining salts compared to

TABLE II. Mean relative error (MRE) between simulation data and experimental results for dielectric constant and molar density of 11 alkali halide brines studied with different ion force fields. All simulations were carried out with the TIP4P/ ε water model and at 298 K, except where noted otherwise.

		10	0 MF	$\operatorname{RE}(\varepsilon)$				
Model	HMN	Fyta	JJ	RH	JC	KBFF	RDVH	MP
salt								
NaCl	5.7	18.9	4.7	6.1	26.8	19.9	24.9	20.5
NaCl (SPC/E)	7.2	25.1	9.7	11.6	33.1	25.8	33.5	25.5
NaCl (278 K)	6.6	13.6	8.9	4.5	20.9	14.6	20.2	16.2
NaCl (308 K)	4.2	20.6	3.2	3.9	26.9	19.9	25.7	20.8
LiCl	6.7		5.3		22.2	22.0	22.0	14.3
RbCl			22.7	24.2	15.9	12.8	13.0	12.4
KCl	11.8	19.5	15.3	17.0	21.0	16.4	18.5	20.2
CsCl	6.7	23.4	27.3	22.9	24.7	25.7	22.9	23.8
LiI	10.9		13.3	7.7	22.4	24.9	18.1	22.3
NaI	6.3	25.4	11.8	12.2	27.4	26.6	26.3	22.5
KI	4.5	33.3	9.9	18.9	27.5	17.8	25.3	25.1
CsI	3.7				21.9	19.5	19.3	19.5
KF	6.6	5.9	18.3	6.3	16.5	17.1	18.6	7.7
CsF	17.5	5.8	24.9	18.4	8.3	20.3	21.8	5.2
Average	7.6	18.6	13.8	12.4	22.5	20.2	22.2	18.3
		10	0 MR	$\operatorname{RE}(\rho)$				
NaCl	2.3	1.1	1.8	0.3	0.2	0.6	1.3	0.2
NaCl (SPC/E)	2.2	1.0	1.8	0.4	0.2	0.6	1.8	0.3
NaCl (278 K)	1.6	0.8	1.3	0.1	0.4	0.3	1.2	0.2
NaCl (308 K)	1.6	1.5	1.2	0.2	0.2	0.9	1.2	0.6
LiCl	1.7		0.8	0.5	2.3	0.7	1.8	1.2
RbCl			5.6	0.4	0.3	1.6	2.9	0.5
KCl	0.1	0.4	1.0	0.1	0.1	0.2	0.7	0.3
CsCl	0.2	0.3	5.7	0.6	0.4	1.2	2.5	0.3
LiI	4.1		2.5	2.0	4.9	3.0	0.8	3.9
NaI	1.8	1.7	1.9	0.7	1.5	3.1	0.4	1.7
KI	0.4	1.0	4.4	1.8	2.0	3.2	1.1	2.4
CsI	0.7	0.5	4.5	1.2	0.7	2.1	0.9	0.7
KF	0.6	1.7	0.4	0.6	2.3	1.3	0.8	1.2
CsF	3.0	1.1	3.1	1.7	3.3	0.2	0.6	2.5
Average	1.6	1.0	2.6	0.8	1.3	1.4	1.3	1.1

the other force fields. The best overall agreement with experimental density data was obtained with the RH-L force field and the worst with the JJ force field.

The predictions of the dielectric constant generally scatter more widely around the experimental data than those of the solution density. For the dielectric constant, at best average deviations of 7.6 % (HMN-S) and at worst deviations of 22.5 % (JC-T) were observed. Dielectric constant calculations with HMN-S were superior to other models for all studied salts. The best representation of the dielectric constant of NaCl brines at all temperatures was given by the RH-L+TIP4P/ ε model combination. The worst agreement for the dielectric constant, with discrepancies to experiments of at least 20 %, was found for the KBFF, RDVH and JC-T force fields, but all of these models yielded much better molar den-

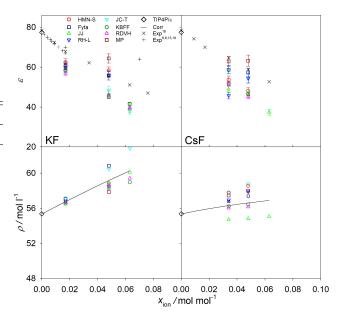


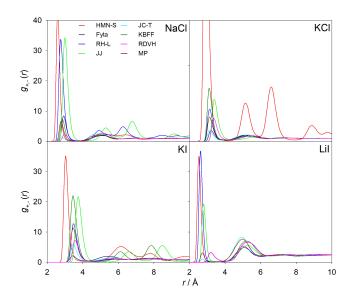
FIG. 5. Dielectric constant (top) and molar density (bottom) over true ion mole fraction of alkali fluoride brines at 298 K and 1 bar. Simulation results for eight ion force fields combined with the TIP4P/ ε water model are represented as open colored symbols as shown in the legend. Experimental data¹⁸ for comparison for KF and CsF brines are represented by cross symbols. Other experimental data for KF brine^{6,9,13,19} are depicted by plus symbols. The prediction for pure water is given by a diamond symbol. A correlation⁸⁷ to experimental density data is shown as a black line. Statistical uncertainties of the molar density are within symbol size.

sity predictions. An overview of the deviations between simulation and experiment is provided in Table II.

In general, low concentration simulations gave better predictions for dielectric constant and molar density, which is not surprising as the water-water interactions dominate the solution behavior. By contrast, the dielectric constant of KCl brine, calculated only up to $0.024 \text{ mol mol}^{-1}$ (as no reliable data for the dielectric constant are available at higher concentrations) was in poor agreement with experimental data.

Qualitative disagreement with experimental dielectric constant data, as well as over- and underestimation of density data, was found for some alkali fluoride models, cf. Figure 5. Taking into account that difficulties in modeling alkali fluorides, especially regarding the solubility, have been reported in the literature,^{37–39,50} we suggest that results for these brines should be analyzed with care.

Interestingly, the models that performed overall best for the dielectric constant (HMN-S, RH-L and JJ) in general led to higher values for the dielectric constant compared to the other models. In some cases, these ion force fields even overestimated the experimental dielectric constant for more concentrated chloride brines, e.g. NaCl or LiCl. Predictions higher than the experimental



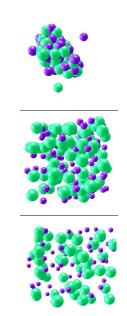


FIG. 6. Cation-anion radial distribution function over distance of aqueous NaCl (0.076 mol mol⁻¹), KCl (0.024 mol mol⁻¹), KI (0.083 mol mol⁻¹) and LiI (0.083 mol mol⁻¹) solutions at 298 K and 1 bar. Simulation results for eight ion force fields combined with the TIP4P/ ε water model are represented by colored lines as shown in the legend.

dielectric constant were also found for those model sets in case of aqueous LiCl before.⁷⁰

Some of the models yielded much higher predictions of the dielectric constant than others, especially for more concentrated brines. In order to find a possible explanation, the cation-anion RDF $g_{+,-}$ (r) of four electrolyte solutions (NaCl, KCl, KI and LiI) was sampled, as shown in Figure 6.

Positions and magnitudes of the first RDF maximum, corresponding to the contact distance and likelihood of contact ion pairs (CIP), respectively, differ significantly between various model sets. The lowest contact distances were observed for the HMN-S and JC-T models and the highest for the JJ and RDVH models. Positions of the first RDF maximum shift to larger distances with increasing cation or anion size.

Very high first peaks, suggesting excessive ion pairing, were formed by the HMN-S model set for the studied salt solutions in Figure 6. In addition, the ordered oscillating behavior of subsequent peaks obtained with this model set (except for LiI) indicates crystal formation in the aqueous solution below the experimental solubility limit. Pronounced oscillations of the RDF curves, which suggest crystallization, were also observed in case of the RH-L and JJ models for NaCl, as well as for the KBFF and JJ models for KI. Models that have the lowest first peaks for NaCl, i.e. JC-T and RDVH, accompanied by MP and KBFF for LiI, gave a higher second than first RDF maximum, which corresponds to a higher likelihood of solvent-shared ion pairs (SIP) than CIP. In SIP, a single solvent shell separates the partner ions, while in CIP

FIG. 7. Snapshots of NaCl configurations modeled by HMN-S (top), MP (center) and RDVH (bottom) force fields in TIP4P/ ε water at 298 K, 1 bar and 0.076 mol mol⁻¹. Violet and green spheres represent Na⁺ and Cl⁻, respectively. Water particles are left out for clarity.

there are no solvent molecules between the cations and anions.⁴ In nearly all other cases, higher first peaks were followed by lower subsequent peaks, as it would be expected for a solution of a completely dissociated salt.

The magnitude of the RDF maximum can provide a relationship with the over- and underestimation of the dielectric constant for a given ion model. In order to establish a connection between the present predictions and excessive ion pairing and solubility, the number of contact ion pairs n^{CIP} was calculated from the cationanion RDF according to Eq. (5), as listed in Table III.

Based on the number of contact ion pairs, the solubility of 1:1 electrolytes in water may be roughly estimated with the ionic pairs rule proposed by Benavides *et al.*⁵² For salts with a model solubility limit below 10 mol kg^{-1} , $n^{\text{CIP}} > 0.5$ is an indicator for a supersaturated solution. i.e. the solubility of the model is lower than the currently investigated concentration. The rule of Benavides et al. is, however, not applicable for salts with a model solubility limit larger than 10 mol kg^{-1} . Salts with an experimental solubility limit above 10 mol kg^{-1} , such as LiCl, CsCl, LiI, NaI and KF, were not treated by the ionic pairs rule. Instead, for these brines, we looked at the snapshot of the final sampled configuration to check visually whether the solution was supersaturated or crystals were formed. All brines and ion force fields studied in this manner at the highest concentration are listed in Table III.

Models with $n^{\text{CIP}} > 0.5$ predict a higher dielectric constant than models with $n^{\text{CIP}} < 0.5$, as less water molecules are needed to form hydration shells and thus

TABLE III. Number of contact ion pairs n^{CIP} calculated from the cation-anion RDF of each brine and each ion force field at 298 K, 1 bar and the highest investigated concentration. The asterisk indicates a supersaturated solution ($n^{\text{CIP}} > 0.5$) according to the ionic pairs rule of Benavides *et al.*⁵², which can be employed for brines with a model solubility limit below 10 mol kg⁻¹. Brines with an experimental solubility limit above 10 mol kg⁻¹ are listed below the horizontal line. For these brines, supersaturated solutions indicated by an asterisk were observed in the snapshots of the final configuration for the investigated model combination.

Model	HMN	Fyta	JJ	RH	JC	KBFF	RDVH	MP
salt								
NaCl	2.97*	0.84^{*}	3.74^{*}	2.97^{*}	0.14	0.69^{*}	0.28	0.57^{*}
$NaCl_{SPC/E}$	2.85^{*}	0.85^{*}	3.69^{*}	2.97^{*}	0.15	0.66^{*}	0.26	0.57^{*}
$NaCl_{278K}$	2.69^{*}	0.78^{*}	3.25^{*}	2.23^{*}	0.10	0.62^{*}	0.27	0.52^{*}
$NaCl_{308K}$	1.13^{*}	0.11	1.42^{*}	1.57^{*}	0.09	0.12	0.24	0.39
RbCl			0.20	0.74^{*}	0.52^{*}	1.61^{*}	1.28^{*}	0.69^{*}
KCl	3.78^{*}	0.22	0.90^{*}	0.55^{*}	0.21	0.80^{*}	0.49	0.43
KI	3.61^{*}	0.45	4.38^{*}	2.33^{*}	1.16^{*}	3.82^{*}	1.46^{*}	2.10^{*}
CsI	4.30^{*}	0.91^{*}	2.43^{*}	0.70^{*}	0.97^{*}	1.24^{*}	1.36^{*}	0.98^{*}
CsF	1.63^{*}	2.20^{*}	0.01	1.63^{*}	0.86^{*}	0.57^{*}	0.83^{*}	1.90^{*}
LiCl	1.36^{*}		1.66^{*}	1.81*	0.32	0.27	0.48	0.64
CsCl	4.26^{*}	1.09	1.59^{*}	1.18	0.94	0.20	1.57	0.86
LiI	0.75^{*}		0.56	0.85^{*}	0.03	0.04	0.25	0.08
NaI	2.65^{*}	0.51	2.00^{*}	1.44^{*}	0.15	0.21	0.19	0.48
\mathbf{KF}	3.10^{*}	1.40^{*}	0.14	2.71^{*}	0.61	0.71	0.73	3.58^{*}

the remaining ones can orient freely.

If crystallization was clearly indicated, as in Figure 6 and Table III for the NaCl models HMN-S, RH-L and JJ, the increase of the dielectric constant was more pronounced. In other words, the reduced model solubility causes premature precipitation, which in turn increases the dielectric constant. Striking differences for the solution structure of a NaCl brine are displayed by snapshots of the microscopic configurations in Figure 7. We also prepared snapshots of the final configuration of each investigated model combination. The visual impression that they provide was always consistent with the rule of Benavides *et al.* Due to the large number of systems, we refrained from including all these snapshots here or in the Supplementary material.

To obtain a compact overview of the connection between the prediction for the dielectric constant and the number of contact ion pairs, ε is shown as a function of n^{CIP} for each brine and each ion force field at the highest investigated concentration in Figures 8 and 9.

For most brines, a clear trend is visible in Figures 8 and 9: points are either found in the lower left or in the upper right part of the plot. This means that all ion force fields that yielded a comparably high dielectric constant at high concentrations were well beyond their solubility limit. The models that did not show this premature precipitation behavior also always predicted a lower dielectric constant.

Hence, a higher number of contact ion pairs caused

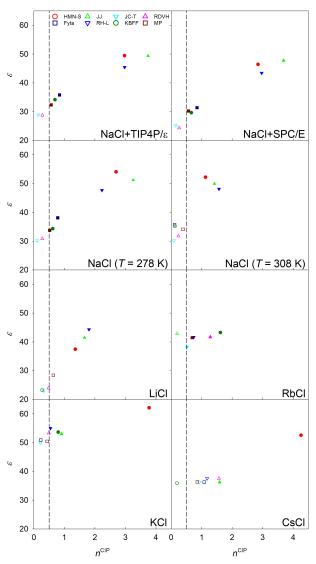


FIG. 8. Dielectric constant as a function of the number of contact ion pairs for alkali chloride brines at the highest investigated concentration. Solid symbols indicate that the solution was likely supersaturated or premature crystallization was observed. The dashed vertical line at $n^{\rm CIP}=0.5$ indicates the limit beyond which this is likely to occur according to the rule of Benavides *et al.*⁵²

an increase of the dielectric constant and thus led to a seemingly better quantitative agreement with experimental data. The argument of a very recent study by Seal *et al.*⁷² that stronger ion pairing leads to a reduced dielectric constant decline at higher NaCl concentrations is thus confirmed here. Another study of the dielectric spectra of NaCl solutions using the GROMOS⁹⁰ force field also found an increment of the dielectric constant due to excessive ion pairing.⁹¹

Table III clearly indicates that most of the ion force fields showed premature precipitation for many of the studied brines. Precipitation not only falsifies results

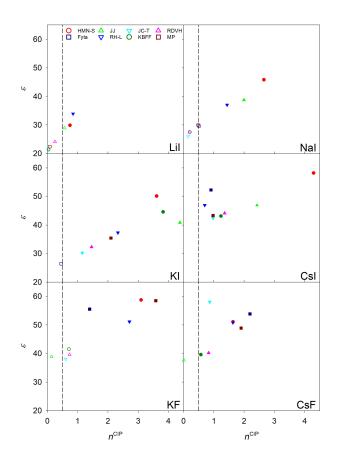


FIG. 9. Dielectric constant as a function of the number of contact ion pairs for alkali iodide and fluoride brines at the highest investigated concentration. Solid symbols indicate that the solution was likely supersaturated or premature crystallization was observed. The dashed vertical line at $n^{\rm CIP} =$ 0.5 indicates the limit beyond which this is likely to occur according to the rule of Benavides *et al.*⁵²

for the dielectric constant, but also for any other thermodynamic property. Exceptions were the JC-T and RDVH force fields, for which only four of the investigated brines were probably above the solubility limit. In fact, Moučka *et al.*⁴⁶ tested the solubility of thirteen force fields for NaCl brine with SPC/E water, among them all ion force fields of this study, except for JJ and MP. They also found that only two ion model sets (JC-SPC/E and the medium model set of Horinek *et al.*) did not show premature precipitation.

The seemingly good reproduction of both dielectric constant and molar density with some ion models is, therefore, probably associated with premature crystallization. Hence, it may be only a coincidence of an unphysical behavior of the respective model, which is thus not useful.

The increase of the dielectric constant with strong ion pairing is contrary to the expected existence of real ion pair species in strong 1:1 electrolyte solutions. The dielectric constant has to be very low ($\varepsilon < 30$) under ambient conditions for uni-univalent ion pairs to be formed.⁴ In fact, the strong drop of the sampled dielectric constant of aqueous NaCl observed before at higher temperatures was the cause of enhanced ion pairing.⁹²

Since strong attraction of cation and anion in solution initiates excessive contact ion pairing, causes for the specific dielectric constant behavior of some models may eventually lie in unbalanced force field parameters. In general, similarly sized ions (small-small and large-large) favor ion association in water, while a combination of small and large ions readily dissociates.^{93–95}

In addition to that is the counterintuitive behavior of the LJ energy parameter for ions in water, for which lower values paradoxically lead to a softer repulsion and stronger electrostatic attraction.^{41,91}

Since all repulsive and dispersive interactions were described in this study with 12 LJ parameters for like and unlike ion and water interactions, it is not straightforward to establish a clear connection between the LJ parameters and the present results. However, in order to relate the findings of this study to the force field parameters of the ions, the ratio between the cation and anion LJ size and energy parameters for each force field is listed in Table S.X.

Ion models that predicted a high dielectric constant, such as HMN-S, had generally similar cation and anion sizes and low LJ energy parameters. On the other hand, ion models with higher LJ energies and strong deviations to the size ratio of unity, such as JC-T, gave a lower number of contact ion pairs. Low LJ cation diameters of the JC-T models are an indication for an increased model solubility⁹⁶ and stronger ion solvation³⁴ in water.

Eleven a queous alkali halide solutions were studied with respect to the ion concentration dependence of the dielectric constant and molar density under ambient conditions by molecular simulation. For that purpose, eight non-polarizable force fields for the ions and the TIP4P/ ε model for water were employed.

Almost all ion force fields predicted a decrease of the dielectric constant with increasing salt concentration for the considered electrolyte solutions. The quantitative agreement with experimental data was poor in many cases. However, the drop of the dielectric constant with increasing temperature for NaCl brine was confirmed. The solution density was predicted well with most ion models. The TIP4P/ ε water model somewhat improved the predictions of the ion force fields regarding dielectric constant and molar density compared to SPC/E.

The key finding of this study is that the generally higher dielectric constant predictions of some ion models are associated with excessive ion pairing and/or premature crystallization. This may explain the seemingly good dielectric constant and density agreement with experiments for these models or in some cases the overestimation of the experimental dielectric constant. The cause probably lies in unbalanced force field parameters and hence reduced model solubility. These findings are contrary to the expected formation of ion pairs with a strong drop of the dielectric constant for 1:1 electrolyte solutions under ambient conditions.

It is clear that further investigations of the dielectric constant of aqueous electrolyte solutions require ion models predicting solubilities close to experimental ones. For example, the recently proposed Madrid-2019 force fields^{97,98} with scaled charges showed good solubility predictions. Reduced charges were used in the excellent reproduction of the experimental dielectric constant for NaCl, NaBr, KCl and KBr solutions with the TIP4P/ $\varepsilon^{88,99}$ water model. However, these findings were compromised with poor solubility and activity predictions of the NaCl/ ε +TIP4P/ ε combination.¹⁰⁰ In the light of the present conclusions, questions arise whether an ion parametrization to the dielectric constant makes sense without an adequate reproduction of the experimental solubility limit.

SUPPLEMENTARY MATERIAL

See Supplementary material for tabulated numerical simulation values and force field parameters used for the water and ion models in this study.

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Supplementary material

Dielectric Constant and Density of Aqueous Alkali Halide Solutions by Molecular Dynamics: A Force Field Assessment

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TABLE S.I. Force field parameters and magnitude of the dipole moment resulting from the distribution of partial charges of the TIP4P/ ε and SPC/E water models.^a

Model	σ (Å)	ε/k_B (K)	$q_{\rm H}$ (e)	$d_{\rm OM}$ (Å)	μ (D)
$\mathrm{TIP4P}/\varepsilon$	3.165	93.0	0.527	0.105	2.434
$\mathrm{SPC/E}$	3.166	78.2	0.424		2.351

^a The negative charge is placed on a fourth (virtual) site M in TIP4P water models so that $q_{\rm M} = -2q_{\rm H}$.

TABLE S.II. Lennard-Jones size and energy parameters of alkali ion force fields considered in this work.

		Li^+		Na^+		K^+		Rb^+		Cs^+	Ref.
Model	σ (Å)	ε/k_B (K))								
HMN-S	2.87	0.07	3.81	0.07	4.53	0.07			5.17	0.07	[1]
Fyta			2.58	50.3	2.69	293.5			3.49	39.1	[2,3]
JJ	2.87	0.25	4.07	0.25	5.17	0.25	5.60	0.25	6.20	0.25	[4]
RH-L	2.68	0.41	3.21	2.15	3.43	23.8	3.62	41.3	3.91	66.6	[5]
JC-T	1.44	52.4	2.18	84.8	2.83	140.7	3.05	218.1	3.36	198.6	[6]
KBFF	1.82	84.2	2.45	38.5	3.34	15.6	3.62	18	4.13	0.78	[7,8]
RDVH	1.88	200	1.89	200	2.77	200	3.26	200	3.58	200	[9,10]
MP	1.72	29	2.50	39.4	3.18	59.6	3.30	121.1	3.44	252.4	[11]

TABLE S.III. Lennard-Jones size and energy parameters of halide ion force fields considered in this work.

		F^{-}		Cl ⁻		Br^-		I^-	Ref.
Model	σ (Å)	ε/k_B (K)	-						
HMN-S	3.43	55.3	4.40	50.5	4.83	25.3	5.33	19.2	[1]
Fyta	4.16	1.80	4.40	50.35	4.63	49.6	5.83	1.80	[2,3]
JJ	3.05	357.5	4.02	357.5	4.28	357.5	4.81	357.5	[4]
RH-L	3.30	47.5	4.09	82.0	4.31	104.5	4.76	114.1	[5]
JC-T	4.52	0.79	4.92	5.87	4.93	15.3	5.26	21	[6]
KBFF	3.70	120.3	4.40	56.5	4.76	36.1	5.35	24.1	[7,8]
RDVH	3.66	200	4.41	200	4.54	200	4.78	200	[9,10]
MP	3.95	3.26	4.61	12.6	4.81	18.1	5.20	21.3	[11]

TABLE S.IV. Experimental and simulation data for dielectric constant and density of pure water with the TIP4P/ ε and SPC/E water models. The numbers in parentheses denote the statistical uncertainties of the dielectric constant in the last given digits, while those of the density are negligibly small.

	T	= 278 K	T =	298 K	T =	308 K
	ε	$\rho \; (\mathrm{mol} \; \mathbf{l}^{-1})$	ε	$\rho \; (\mathrm{mol} \; \mathbf{l}^{-1})$	ε	$\rho \pmod{l^{-1}}$
$Exp.^{12,13}$		55.51	78.5	55.35	74.9	55.35
$TIP4P/\varepsilon$	85(2)		77.4(15)	55.36	74.3(12)	55.18
SPC/E			69.0(13)	55.26		

TABLE S.V. Experimental and simulation data for the dielectric constant ε of alkali chloride brines at 298 K, 1 bar and specified concentrations. The asterisk indicates a simulation point where a supersaturated solution is expected according to the ionic pairs rule of Benavides *et al.*¹⁴ or for which the snapshot of the final sampled configuration clearly showed the presence of a crystal phase. The numbers in parentheses denote the statistical uncertainties in the last given digits.

salt	NaC	Cl + TIP4	P/ε	Na	Cl + SPC	/E
$x_{\rm ion} \; ({\rm mol} \; {\rm mol}^{-1})$	0.026	0.049	0.076	0.026	0.049	0.076
$\operatorname{Exp.}^{13}$	62.2	52.4	43.5	62.2	52.4	43.5
Model						
HMN-S	60.3(13)	56.8 (12)	*	54.6 (11)	51.3(1)	*
Fyta	51.9(11)	41.3(8)	*	47.4(9)	38.6(7)	*
JJ	56.9(12)	52.9(11)	*	52.1(10)	50.7(11)	*
RH-L		49.6(11)	*	50.9(9)	49.8(10)	*
JC-T	50.3(11)	· · ·	28.9(6)	46.4(7)	35.7~(6)	25.2(4)
KBFF	49.6(9)	42.2(9)	*	47.2(8)	37.9(7)	*
RDVH	52.4(12)		28.7(5)	47.9(10)		24.3(4)
MP	51.3(7)	40.1(6)	*	47.1(9)	38.4(6)	*
salt		LiCl			RbCl	
$x_{\rm ion} \ ({\rm mol} \ {\rm mol}^{-1})$	0.025	0.049	0.095	0.022	0.057	0.063
$Exp.^{15,16}$	60.3	48.6	33.3	66.0	55.0	53.7
Model						
HMN-S	52.7(11)	48.2 (12)	*			
Fyta						
JJ	· · ·	50.7(11)	*		42.8(10)	· · ·
RH-L		51.0(11)	*	54.7(12)		*
JC-T	49.3(12)		23.0(5)	54.6(11)		*
KBFF	50.7(10)		23.2(5)	54.7(12)	. ,	*
RDVH	51.2(10)		23.9(4)	56.2(11)		*
MP	51.7 (11)	41.9(1)	28.4(6)	57.7 (13)	41.4(8)	*
salt		KCl			CsCl	
$x_{\rm ion} \; ({\rm mol} \; {\rm mol}^{-1})$	0.007	0.015	0.024	0.013	0.023	0.068
Exp. ^{16,17}	75.5	71.2	65.9	72.0	68.0	57.6
Model						
HMN-S	66.0(14)	64.2(15)	*	67.2(17)	63.4(15)	*
Fyta		58.3(12)	50.9(10)		56.3(12)	
JJ		61.8(13)	*		52.6(12)	
RH-L	64.6(14)	58.2(12)	*	62.8(13)	53.6(12)	37.6(9)
JC-T		57.7(12)			53.4(12)	
KBFF		59.0(13)	*		53.0(11)	
RDVH		58.1(11)			54.3(12)	
MP	63.3(14)	58.3(12)	50.4(10)	60.8(12)	55.2(11)	36.4(8)

TABLE S.VI. Experimental and simulation data for the dielectric constant ε of alkali iodide and fluoride brines at 298 K, 1 bar and specified concentrations. The asterisk indicates a simulation point where a supersaturated solution is expected according to the ionic pairs rule of Benavides *et al.*¹⁴ or for which the snapshot of the final sampled configuration clearly showed the presence of a crystal phase. The numbers in parentheses denote the statistical uncertainties in the last given digits.

salt		LiI			NaI	
$x_{\rm ion} \; ({\rm mol} \; {\rm mol}^{-1})$	0.034	0.049	0.083	0.013	0.036	0.072
Exp. ^{18–20}	52.4	44.7	32.7	67.0	55.5	45.7
Model						
HMN-S	46.6(11)	39.9(12)	*	· · ·	53.8(11)	*
Fyta		()	.1.	55.9(12)		29.6(5)
JJ	45.9 (8)	38.3 (10)	*		48.4 (10)	*
RH-L	47.0 (12)	· · ·	*	59.5(13)		*
JC-T	43.9 (12)		21.7(4)	57.9 (13)		26.0(5)
KBFF	43.0 (10)	· · ·	21.4(5)	56.0(10)	· · ·	27.5(5)
RDVH	46.2(10)	· · ·	24.1(4)	57.6(12)	· · ·	27.5(5)
MP	44.0 (10)	36.0 (6)	22.4(4)	58.1(12)	44.4(9)	30.0(5)
salt		KI			CsI	
$x_{\rm ion} \; ({\rm mol} \; {\rm mol}^{-1})$	0.034	0.049	0.083	0.017	0.034	0.045
$Exp.^{21,22}$	65.3	55.2	42.6	69.0	62.5	57
Model						
HMN-S	60.3(14)	54.4(13)	*	65.5(15)	61.1(14)	*
Fyta	44.7 (9)	38.5(8)	26.4(5)	58.2(11)	47.3 (11)	*
JJ	55.1(12)	52.9(12)	*	58.9(14)	48.8 (11)	*
RH-L	51.4(12)	46.1 (9)	*	59.3(13)	51.1 (11)	*
JC-T	46.6(9)	40.7(8)	*	57.4(12)	45.6(9)	*
KBFF	50.6(11)	47.9 (10)	*	56.2(12)	49.7 (11)	*
RDVH	47.4 (10)	42.4 (8)	*	57.4(13)	48.8 (10)	*
MP	48.8 (11)	41.5 (7)	*	57.2 (12)	48.8 (11)	*
salt		KF			CsF	
$x_{\rm ion} \; ({\rm mol} \; {\rm mol}^{-1})$	0.017	0.048	0.063	0.034	0.048	0.063
$\operatorname{Exp.}^{23}$	67.9	55.4	51.1	64.6	58.8	52.7
Model						
HMN-S	62.5(17)	58.4 (17)	*	53.7(16)	48.2 (15)	*
Fyta	· · ·	55.8 (23)	*		57.3 (22)	*
JJ	· · ·	46.2 (18)	38.9(19)	. ,	. ,	37.4(16)
RH-L	· · ·	55.8 (22)	*	()	54.3 (23)	*
JC-T	· · ·	48.5 (20)	38.1 (17)	· · · ·	54.7 (24)	*
KBFF	· · ·	45.1 (10)	· · ·	51.8 (11)		*
RDVH	56.9 (11)	. ,	39.6(8)	. ,	45.1 (10)	*
MP	· · ·	64.4 (23)	*	()	63.3 (28)	*

TABLE S.VII. Correlation and simulation data for the molar density ρ in mol l^{-1} of alkali halide brines at 298 K, 1 bar and specified concentrations. The asterisk indicates a simulation point where a supersaturated solution is expected according to the ionic pairs rule of Benavides *et al.*¹⁴ or for which the snapshot of the final sampled configuration clearly showed the presence of a crystal phase. Statistical uncertainties are negligibly small.

·	· · · · · · · · · · · · · · · · · · ·			
salt	$NaCl + TIP4P/\varepsilon$	NaCl + SPC/E	LiCl	RbCl
$x_{\rm ion} \; ({\rm mol} \; {\rm mol}^{-1})$	0.026 0.049 0.076	$0.026 \ 0.049 \ 0.076$	$0.025 \ 0.049 \ 0.095$	$0.022 \ 0.057 \ 0.063$
Corr. ²⁴	56.71 57.87 59.16	56.71 57.87 59.16	56.69 57.96 60.34	56.71 57.87 59.16
Model				
HMN-S	55.90 56.01 *	55.97 56.09 *	56.05 56.70 *	
Fyta	56.33 56.99 *	56.41 57.02 *		
JJ	56.13 56.40 *	56.15 56.40 *	56.48 57.21 *	53.94 51.85 51.54
RH-L	56.84 57.70 *	56.90 57.60 *	56.55 57.55 *	55.72 55.83 *
JC-T	56.89 58.05 59.23	56.93 58.03 59.16	56.09 56.76 58.14	55.50 55.33 *
KBFF	56.55 57.38 *	56.58 57.34 *	57.09 58.51 60.57	55.04 54.26 *
RDVH	56.44 57.15 57.82	56.32 56.85 57.35	56.43 57.19 58.17	54.61 53.25 *
MP	56.87 57.95 *	56.96 58.00 *	56.98 57.50 58.99	55.57 55.13 *
salt	KCl	CsCl	LiI	NaI
$x_{\rm ion} \; ({\rm mol} \; {\rm mol}^{-1})$	0.007 0.015 0.024	0.013 0.023 0.068	0.034 0.049 0.083	0.013 0.036 0.072
Corr. ²⁴	55.49 55.66 55.93	55.07 54.89 54.29	55.49 55.66 55.93	55.07 54.89 54.29
Model				
HMN-S	55.48 55.55 *	55.16 55.01 *	53.38 52.67 *	54.77 53.72 *
Fyta	55.60 55.76 56.04	54.91 55.10 54.18		55.02 54.37 53.44
JJ	55.18 54.91 *	53.83 52.70 48.45	54.22 53.59 *	54.74 53.67 *
RH-L	55.57 55.73 *	55.07 54.78 53.36	54.43 53.96 *	55.12 54.59 *
JC-T	55.56 55.69 55.88	55.18 54.99 53.89	53.53 52.83 51.48	$55.08 \ 54.46 \ 53.51$
KBFF	55.45 55.45 *	54.93 54.56 52.75	54.33 53.86 52.73	54.75 53.63 52.04
RDVH	55.40 55.31 55.19	54.64 54.05 51.48	55.18 54.99 54.44	55.36 55.15 54.62
MP	55.75 55.63 55.76	54.91 55.10 54.18	53.99 53.37 52.02	55.03 54.34 53.34
salt	KI	CsI	KF	CsF
$x_{\rm ion} \ ({\rm mol} \ {\rm mol}^{-1})$	0.034 0.049 0.083	$0.017 \ 0.034 \ 0.045$	$0.017 \ 0.048 \ 0.063$	0.034 0.048 0.063
$\operatorname{Corr.}^{24}$	54.12 53.72 53.02	53.87 52.83 52.31	56.75 59.17 60.28	56.31 56.61 56.89
Model				
HMN-S	53.91 53.46 *	54.20 53.25 *	57.09 58.85 *	57.78 58.57 *
Fyta	53.95 53.34 51.99	54.10 53.17 *	57.10 60.85 *	56.77 57.39 *
JJ	52.15 50.92 *	52.19 49.69 *	56.47 58.90 60.12	54.71 54.83 55.05
RH-L	53.36 52.53 *	53.56 51.89 *	56.75 58.48 *	56.96 57.89 *
JC-T	53.28 52.42 *	53.73 52.21 *	57.10 60.46 62.83	57.81 58.79 *
KBFF	52.61 51.82 *	53.22 51.27 *	56.55 58.29 58.97	56.22 56.41 *
RDVH	53.71 52.98 *	53.66 52.07 *	56.72 58.67 59.43	56.03 56.14 *
MP	53.08 52.17 *	53.70 52.24 *	56.53 57.98 *	57.45 58.00 *
	55.00 52.11	55.10 52.24	00.00 01.00	01.10 00.00

TABLE S.VIII. Experimental and simulation data for the dielectric constant ε and molar density ρ of sodium chloride brines at 278 and 308 K, 1 bar and specified concentrations. The asterisk indicates a simulation point where a supersaturated solution is expected according to the ionic pairs rule of Benavides *et al.*¹⁴ or for which the snapshot of the final sampled configuration clearly showed the presence of a crystal phase. The numbers in parentheses denote the statistical uncertainties of the dielectric constant in the last given digits, while those of the density are negligibly small.

salt	Na	$\operatorname{Cl}(T=2)$	78 K)	Na	Cl $(T = 30)$	08 K)
$x_{\rm ion} \; ({\rm mol} \; {\rm mol}^{-1})$	0.026	0.049	0.078	0.026	0.049	0.064
	Diel	ectric cons	stant ε	Diele	ectric cons	tant ε
Exp. ¹³	66.4	52.8	40.1	59.5	50.9	46.8
Model						
HMN-S	63.1 (14) 57.2 (9)	*	57.9(9)	53.8(7)	*
Fyta		46.2 (11)) *	49.5 (9)	39.9 (6)	35.7(6)
JJ) 58.6 (12)		55.8(9)	51.0(7)	*
RH-L	60.8 (18) 52.5 (16)) *	55.0(12)	51.0(11)	*
JC-T	55.5(15)) 41.4 (14)) 30.2 (12)	48.8 (10)	37.1 (10)	30.2(10)
KBFF	56.4(15)) 45.3 (12)) *	50.5(9)	40.9(7)	35.2(6)
RDVH	55.7 (15) 41.5 (11)) 30.8 (9)	48.9(9)		31.8(6)
MP	56.4(17)) 43.6 (17)) *	50.5(10)	40.6(10)	34.1(10)
salt	Na	Cl (T = 2)	78 K)	Na	Cl $(T = 30)$	08 K)
$x_{\rm ion} \; ({\rm mol} \; {\rm mol}^{-1})$	0.026	0.049	0.078	0.026	0.049	0.064
		$\rho~({\rm mol}~{\rm l}^-$	1)		$\rho \pmod{l^{-1}}$)
$\operatorname{Exp.}^{13}$, $\operatorname{Corr.}^{24}$	57.09	58.29	59.67	56.52	57.67	58.27
Model						
HMN-S	56.19	56.35	*	55.68	55.80	*
Fyta	56.67	57.39	*	56.11	56.73	57.07
JJ	56.42	56.75	*	55.90	56.18	*
RH-L	57.17	58.27	*	56.58	57.33	*
JC-T	57.30	58.54	59.89	56.64	57.75	58.41
KBFF	56.94	57.85	*	56.31	57.11	57.54
RDVH	56.84	57.64	58.41	56.20	56.88	57.23
MP	56.96	58.00	*	56.84	57.85	58.73

TABLE S.IX. Experimental and simulation data for the dielectric constant ε and molar density ρ in mol l⁻¹ of alkali halide brines at 298 K, 1 bar and the highest concentration. The asterisk indicates a simulation point where a supersaturated solution is expected according to the ionic pairs rule of Benavides *et al.*¹⁴ or for which the snapshot of the final sampled configuration clearly showed the presence of a crystal phase. Simulation results for solutions for which this is not the case are provided in Tables S.V to S.VIII. The numbers in parentheses denote the statistical uncertainties of the dielectric constant in the last given digits, while those of the density are negligibly small.

14	N ₂ Ol	N ₂ Ol	N ₂ Ol	M-CI	T:CI	DL CI	VOL
salt	NaCl	NaCl $_{\rm SPC/E}$				RbCl	KCl
$x_{\text{ion}} \pmod{\text{mol}^{-1}}$	0.076	0.076	0.078	0.064	0.095	0.063	0.024
	ε	ε	ε	ε	ε	ε	ε
$Exp.^{13,15-17}$	43.5	43.5	40.10	46.80	33.3	53.7	65.9
Model							
HMN-S	49.5 (11)*	46.4 (10)*	54.1 (7)*	52.2 (5)*	37.4 (10)*		62.1 (13)*
Fyta	$35.8 (7)^*$		$38.1 (8)^*$				
JJ		47.7 (10)*	$51.2 (9)^*$	49.9 (7)*	41.5 (11)*	B (A (A) *	$53.0(11)^*$
RH-L JC-T	$45.4 (11)^*$	$43.4 (9)^*$	47.7 (14)*	48.1 (11)*	44.3 (11)*	$36.9(8)^*$ $35.4(7)^*$	$55.0 (13)^*$
KBFF	34.1 (7)*	$29.6 (5)^*$	34.4 (9)*			$41.7 (8)^*$	53.6 (10)*
RDVH	0.1.2 (1)	(0)	0 (0)			$37.3 (7)^*$	
MP	$32.3 (11)^*$	$30.2 (5)^*$	$33.8 (16)^*$			$36.1(7)^*$	
	ρ	ρ	ρ	ho	ho	ρ	ho
$\operatorname{Exp.}^{13}$, $\operatorname{Corr.}^{24}$	59.16	59.16	59.67	58.27	60.34	59.16	55.93
HMN-S	56.24*	56.23*	56.66*	55.94*	58.15*		55.65*
Fyta	57.58*	57.56*	58.08*	F A 0.14	FO 11 *		F 4 4 0 ¥
JJ RH-L	56.83* 58.54*	56.79* 58.50*	57.27* 59.21*	56.34* 57.85*	58.41* 59.43*	55.83*	54.42* 55.94*
JC-T	38.34	38.30	39.21	37.83	09.45 ⁺	55.31*	55.94
KBFF	58.17^{*}	58.02*	58.78^{*}			54.15^*	55.43*
RDVH						53.04*	
MP	59.01*	59.01*	59.01*			55.13*	
salt	CsCl	LiI	NaI	KI	CsI	KF	CsF
$x_{\rm ion} \; ({\rm mol} \; {\rm mol}^{-1})$	0.068	0.024	0.072	0.083	0.045	0.063	0.063
	ε	ε	ε	ε	ε	ε	ε
$\operatorname{Exp.}^{17-23}$	57.6	55.93	45.7	42.6	57	51.1	52.7
Model	-						
HMN-S	53.5 (12)*	29.9 (9)*	45.8 (11)*	50.1 (12)*	58.2 (13)*	58.8 (18)*	51.1 (17)*
HMN-S Fyta	53.5 (12)*					58.8 (18)* 55.5 (26)*	
HMN-S Fyta JJ	53.5 (12)*	28.9 (10)*	38.8 (7)*	40.8 (9)*	$52.2 (11)^*$ 46.9 (12)*	55.5 (26)*	53.9 (24)*
HMN-S Fyta JJ RH-L	53.5 (12)*			40.8 (9)* 37.3 (8)*	$52.2 (11)^* 46.9 (12)^* 46.9 (11)^*$		53.9 (24)* 50.8 (25)*
HMN-S Fyta JJ RH-L JC-T	53.5 (12)*	28.9 (10)*	38.8 (7)*	40.8 (9)* 37.3 (8)* 30.2 (6)*	$52.2 (11)^{*} 46.9 (12)^{*} 46.9 (11)^{*} 42.4 (8)^{*}$	55.5 (26)*	53.9 (24)* 50.8 (25)* 58.1 (29)*
HMN-S Fyta JJ RH-L JC-T KBFF	53.5 (12)*	28.9 (10)*	38.8 (7)*	$\begin{array}{c} 40.8 \ (9)^* \\ 37.3 \ (8)^* \\ 30.2 \ (6)^* \\ 44.6 \ (9)^* \end{array}$	$52.2 (11)^{*} 46.9 (12)^{*} 46.9 (11)^{*} 42.4 (8)^{*} 43.1 (9)^{*}$	55.5 (26)*	53.9 (24)* 50.8 (25)* 58.1 (29)* 39.3 (8)*
HMN-S Fyta JJ RH-L JC-T	53.5 (12)*	28.9 (10)*	38.8 (7)*	40.8 (9)* 37.3 (8)* 30.2 (6)* 44.6 (9)* 32.3 (6)*	$52.2 (11)^{*} 46.9 (12)^{*} 46.9 (11)^{*} 42.4 (8)^{*} 43.1 (9)^{*} 44.1 (8)^{*}$	55.5 (26)* 51.2 (19)*	53.9 (24)* 50.8 (25)* 58.1 (29)* 39.3 (8)* 40.0 (8)*
HMN-S Fyta JJ RH-L JC-T KBFF RDVH		28.9 (10)* 34.0 (11)*	38.8 (7)* 37.0 (7)*	$\begin{array}{c} 40.8 \ (9)^{*} \\ 37.3 \ (8)^{*} \\ 30.2 \ (6)^{*} \\ 44.6 \ (9)^{*} \\ 32.3 \ (6)^{*} \\ 35.4 \ (7)^{*} \end{array}$	$52.2 (11)^{*} 46.9 (12)^{*} 46.9 (11)^{*} 42.4 (8)^{*} 43.1 (9)^{*} 44.1 (8)^{*} 43.3 (9)^{*} $	55.5 (26)* 51.2 (19)* 58.5 (27)*	53.9 (24)* 50.8 (25)* 58.1 (29)* 39.3 (8)* 40.0 (8)* 48.9 (23)*
HMN-S Fyta JJ RH-L JC-T KBFF RDVH MP	ρ	28.9 (10)* 34.0 (11)* ρ	38.8 (7)* 37.0 (7)* ρ	$\begin{array}{c} 40.8 \ (9)^* \\ 37.3 \ (8)^* \\ 30.2 \ (6)^* \\ 44.6 \ (9)^* \\ 32.3 \ (6)^* \\ 35.4 \ (7)^* \end{array}$	$\begin{array}{c} 52.2\ (11)^{\ast}\\ 46.9\ (12)^{\ast}\\ 46.9\ (11)^{\ast}\\ 42.4\ (8)^{\ast}\\ 43.1\ (9)^{\ast}\\ 44.1\ (8)^{\ast}\\ 43.3\ (9)^{\ast}\\ \end{array}$	55.5 (26)* 51.2 (19)* 58.5 (27)* ρ	$53.9 (24)^{*}$ $50.8 (25)^{*}$ $58.1 (29)^{*}$ $39.3 (8)^{*}$ $40.0 (8)^{*}$ $48.9 (23)^{*}$ ρ
HMN-S Fyta JJ RH-L JC-T KBFF RDVH MP	$\frac{\rho}{54.29}$	$\begin{array}{c} 28.9 \ (10)^{*} \\ 34.0 \ (11)^{*} \\ \hline \\ \rho \\ 55.93 \end{array}$	38.8 (7)* 37.0 (7)* ρ 54.29	$\begin{array}{c} 40.8 \ (9)^* \\ 37.3 \ (8)^* \\ 30.2 \ (6)^* \\ 44.6 \ (9)^* \\ 32.3 \ (6)^* \\ 35.4 \ (7)^* \\ \hline \rho \\ 53.02 \\ \end{array}$	$\begin{array}{c} 52.2\ (11)^{\ast}\\ 46.9\ (12)^{\ast}\\ 46.9\ (11)^{\ast}\\ 42.4\ (8)^{\ast}\\ 43.1\ (9)^{\ast}\\ 44.1\ (8)^{\ast}\\ 43.3\ (9)^{\ast}\\ \hline \rho\\ 52.31\\ \end{array}$	55.5 $(26)^*$ 51.2 $(19)^*$ 58.5 $(27)^*$ ρ 60.28	$\begin{array}{c} 53.9 \ (24)^{*} \\ 50.8 \ (25)^{*} \\ 58.1 \ (29)^{*} \\ 39.3 \ (8)^{*} \\ 40.0 \ (8)^{*} \\ 48.9 \ (23)^{*} \\ \hline \rho \\ 57.30 \end{array}$
HMN-S Fyta JJ RH-L JC-T KBFF RDVH MP Corr. ²⁴ HMN-S	ρ	28.9 (10)* 34.0 (11)* ρ	38.8 (7)* 37.0 (7)* ρ	$\begin{array}{c} 40.8 \ (9)^* \\ 37.3 \ (8)^* \\ 30.2 \ (6)^* \\ 44.6 \ (9)^* \\ 32.3 \ (6)^* \\ 35.4 \ (7)^* \end{array}$	$\begin{array}{c} 52.2\ (11)^{*}\\ 46.9\ (12)^{*}\\ 46.9\ (11)^{*}\\ 42.4\ (8)^{*}\\ 43.1\ (9)^{*}\\ 44.1\ (8)^{*}\\ 43.3\ (9)^{*}\\ \hline \rho\\ 52.31\\ \hline 52.55^{*}\\ \end{array}$	55.5 (26)* 51.2 (19)* 58.5 (27)* ρ 60.28 60.29*	$\begin{array}{c} 53.9\ (24)^{*}\\ 50.8\ (25)^{*}\\ 58.1\ (29)^{*}\\ 39.3\ (8)^{*}\\ 40.0\ (8)^{*}\\ 48.9\ (23)^{*}\\ \hline \rho\\ 57.30\\ 59.43^{*}\\ \end{array}$
HMN-S Fyta JJ RH-L JC-T KBFF RDVH MP Corr. ²⁴ HMN-S Fyta	$\frac{\rho}{54.29}$	$\begin{array}{c} 28.9 \ (10)^{*} \\ 34.0 \ (11)^{*} \\ \hline \\ 55.93 \\ 51.42^{*} \end{array}$	38.8 (7)* 37.0 (7)* ρ 54.29	$\begin{array}{c} 40.8 \ (9)^* \\ 37.3 \ (8)^* \\ 30.2 \ (6)^* \\ 44.6 \ (9)^* \\ 32.3 \ (6)^* \\ 35.4 \ (7)^* \\ \hline \rho \\ 53.02 \\ 52.66^* \end{array}$	$\begin{array}{c} 52.2\ (11)^{*}\\ 46.9\ (12)^{*}\\ 46.9\ (11)^{*}\\ 42.4\ (8)^{*}\\ 43.1\ (9)^{*}\\ 44.1\ (8)^{*}\\ 43.3\ (9)^{*}\\ \hline \rho\\ 52.31\\ \hline 52.55^{*}\\ 52.64^{*}\\ \end{array}$	55.5 $(26)^*$ 51.2 $(19)^*$ 58.5 $(27)^*$ ρ 60.28	$\begin{array}{c} 53.9 \ (24)^{*} \\ 50.8 \ (25)^{*} \\ 58.1 \ (29)^{*} \\ 39.3 \ (8)^{*} \\ 40.0 \ (8)^{*} \\ 48.9 \ (23)^{*} \\ \hline \rho \\ 57.30 \end{array}$
HMN-S Fyta JJ RH-L JC-T KBFF RDVH MP Corr. ²⁴ HMN-S	$\frac{\rho}{54.29}$	$\begin{array}{c} 28.9 \ (10)^{*} \\ 34.0 \ (11)^{*} \\ \hline \\ \rho \\ 55.93 \end{array}$	$\begin{array}{c} 38.8 \ (7)^{*} \\ 37.0 \ (7)^{*} \\ \hline \\ 54.29 \\ 52.18^{*} \end{array}$	$\begin{array}{c} 40.8 \ (9)^* \\ 37.3 \ (8)^* \\ 30.2 \ (6)^* \\ 44.6 \ (9)^* \\ 32.3 \ (6)^* \\ 35.4 \ (7)^* \\ \hline \rho \\ 53.02 \\ \end{array}$	$\begin{array}{c} 52.2\ (11)^{*}\\ 46.9\ (12)^{*}\\ 46.9\ (11)^{*}\\ 42.4\ (8)^{*}\\ 43.1\ (9)^{*}\\ 44.1\ (8)^{*}\\ 43.3\ (9)^{*}\\ \hline \rho\\ 52.31\\ \hline 52.55^{*}\\ \end{array}$	55.5 (26)* 51.2 (19)* 58.5 (27)* ρ 60.28 60.29*	$\begin{array}{c} 53.9\ (24)^{*}\\ 50.8\ (25)^{*}\\ 58.1\ (29)^{*}\\ 39.3\ (8)^{*}\\ 40.0\ (8)^{*}\\ 48.9\ (23)^{*}\\ \hline \rho\\ 57.30\\ 59.43^{*}\\ \end{array}$
HMN-S Fyta JJ RH-L JC-T KBFF RDVH MP Corr. ²⁴ HMN-S Fyta JJ RH-L	$\frac{\rho}{54.29}$	$\begin{array}{c} 28.9 \ (10)^{*} \\ 34.0 \ (11)^{*} \\ \hline \\ 55.93 \\ 51.42^{*} \\ 52.02^{*} \end{array}$	$\begin{array}{c} 38.8 \ (7)^{*} \\ 37.0 \ (7)^{*} \\ \hline \\ 54.29 \\ 52.18^{*} \\ 52.04^{*} \end{array}$	$\begin{array}{c} 40.8 & (9)^{*} \\ 37.3 & (8)^{*} \\ 30.2 & (6)^{*} \\ 44.6 & (9)^{*} \\ 32.3 & (6)^{*} \\ 35.4 & (7)^{*} \\ \hline \\ \hline \\ \hline \\ 53.02 \\ 52.66^{*} \\ 48.32^{*} \\ 50.76^{*} \\ \end{array}$	$\begin{array}{c} 52.2\ (11)^{*}\\ 46.9\ (12)^{*}\\ 46.9\ (11)^{*}\\ 42.4\ (8)^{*}\\ 43.1\ (9)^{*}\\ 43.3\ (9)^{*}\\ \hline \\ \hline \\ 52.31\\ \hline \\ 52.55^{*}\\ 52.64^{*}\\ 48.16^{*}\\ 50.87^{*}\\ \end{array}$	$55.5 (26)^{*}$ $51.2 (19)^{*}$ $58.5 (27)^{*}$ ρ 60.28 60.29^{*} 62.95^{*}	$53.9 (24)^{*}$ $50.8 (25)^{*}$ $58.1 (29)^{*}$ $39.3 (8)^{*}$ $40.0 (8)^{*}$ $48.9 (23)^{*}$ ρ 57.30 59.43^{*} 58.76^{*}
HMN-S Fyta JJ RH-L JC-T KBFF RDVH MP Corr. ²⁴ HMN-S Fyta JJ RH-L JC-T	$\frac{\rho}{54.29}$	$\begin{array}{c} 28.9 \ (10)^{*} \\ 34.0 \ (11)^{*} \\ \hline \\ 55.93 \\ 51.42^{*} \\ 52.02^{*} \end{array}$	$\begin{array}{c} 38.8 \ (7)^{*} \\ 37.0 \ (7)^{*} \\ \hline \\ 54.29 \\ 52.18^{*} \\ 52.04^{*} \end{array}$	$\begin{array}{c} 40.8 & (9)^{*} \\ 37.3 & (8)^{*} \\ 30.2 & (6)^{*} \\ 44.6 & (9)^{*} \\ 32.3 & (6)^{*} \\ 35.4 & (7)^{*} \\ \hline \\ \hline \\ \hline \\ 53.02 \\ 52.66^{*} \\ 48.32^{*} \\ 50.76^{*} \\ 50.63^{*} \\ \end{array}$	$\begin{array}{c} 52.2\ (11)^{*}\\ 46.9\ (12)^{*}\\ 46.9\ (11)^{*}\\ 42.4\ (8)^{*}\\ 43.1\ (9)^{*}\\ 44.1\ (8)^{*}\\ 43.3\ (9)^{*}\\ \hline \rho\\ 52.31\\ \hline 52.55^{*}\\ 52.64^{*}\\ 48.16^{*}\\ \end{array}$	$55.5 (26)^{*}$ $51.2 (19)^{*}$ $58.5 (27)^{*}$ ρ 60.28 60.29^{*} 62.95^{*}	$53.9 (24)^{*}$ $50.8 (25)^{*}$ $58.1 (29)^{*}$ $39.3 (8)^{*}$ $40.0 (8)^{*}$ $48.9 (23)^{*}$ $\overline{\rho}$ 57.30 59.43^{*} 58.76^{*} 58.87^{*} 60.49^{*}
HMN-S Fyta JJ RH-L JC-T KBFF RDVH MP Corr. ²⁴ HMN-S Fyta JJ RH-L	$\frac{\rho}{54.29}$	$\begin{array}{c} 28.9 \ (10)^{*} \\ 34.0 \ (11)^{*} \\ \hline \\ 55.93 \\ 51.42^{*} \\ 52.02^{*} \end{array}$	$\begin{array}{c} 38.8 \ (7)^{*} \\ 37.0 \ (7)^{*} \\ \hline \\ 54.29 \\ 52.18^{*} \\ 52.04^{*} \end{array}$	$\begin{array}{c} 40.8 & (9)^{*} \\ 37.3 & (8)^{*} \\ 30.2 & (6)^{*} \\ 44.6 & (9)^{*} \\ 32.3 & (6)^{*} \\ 35.4 & (7)^{*} \\ \hline \\ \hline \\ \hline \\ 53.02 \\ 52.66^{*} \\ 48.32^{*} \\ 50.76^{*} \\ \end{array}$	$\begin{array}{c} 52.2\ (11)^{*}\\ 46.9\ (12)^{*}\\ 46.9\ (11)^{*}\\ 42.4\ (8)^{*}\\ 43.1\ (9)^{*}\\ 44.1\ (8)^{*}\\ 43.3\ (9)^{*}\\ \hline \\ \hline \\ 52.31\\ 52.55^{*}\\ 52.64^{*}\\ 48.16^{*}\\ 50.87^{*}\\ 51.27^{*}\\ \end{array}$	$55.5 (26)^{*}$ $51.2 (19)^{*}$ $58.5 (27)^{*}$ ρ 60.28 60.29^{*} 62.95^{*}	$53.9 (24)^{*}$ $50.8 (25)^{*}$ $58.1 (29)^{*}$ $39.3 (8)^{*}$ $40.0 (8)^{*}$ $48.9 (23)^{*}$ $\overline{\rho}$ 57.30 59.43^{*} 58.76^{*} 58.87^{*}
HMN-S Fyta JJ RH-L JC-T KBFF RDVH MP Corr. ²⁴ HMN-S Fyta JJ RH-L JC-T KBFF	$\frac{\rho}{54.29}$	$\begin{array}{c} 28.9 \ (10)^{*} \\ 34.0 \ (11)^{*} \\ \hline \\ 55.93 \\ 51.42^{*} \\ 52.02^{*} \end{array}$	$\begin{array}{c} 38.8 \ (7)^{*} \\ 37.0 \ (7)^{*} \\ \hline \\ 54.29 \\ 52.18^{*} \\ 52.04^{*} \end{array}$	$\begin{array}{c} 40.8 & (9)^* \\ 37.3 & (8)^* \\ 30.2 & (6)^* \\ 44.6 & (9)^* \\ 32.3 & (6)^* \\ 35.4 & (7)^* \\ \hline \\ \hline \\ \hline \\ 53.02 \\ 52.66^* \\ 48.32^* \\ 50.76^* \\ 50.63^* \\ 50.36^* \\ \end{array}$	$\begin{array}{c} 52.2\ (11)^{*}\\ 46.9\ (12)^{*}\\ 46.9\ (11)^{*}\\ 42.4\ (8)^{*}\\ 43.1\ (9)^{*}\\ 43.1\ (9)^{*}\\ 43.3\ (9)^{*}\\ \hline \\ \hline \\ 52.31\\ 52.55^{*}\\ 52.64^{*}\\ 48.16^{*}\\ 50.87^{*}\\ 51.27^{*}\\ 50.10^{*}\\ \end{array}$	$55.5 (26)^{*}$ $51.2 (19)^{*}$ $58.5 (27)^{*}$ ρ 60.28 60.29^{*} 62.95^{*}	$53.9 (24)^{*}$ $50.8 (25)^{*}$ $58.1 (29)^{*}$ $39.3 (8)^{*}$ $40.0 (8)^{*}$ $48.9 (23)^{*}$ ρ 57.30 59.43^{*} 58.76^{*} 58.87^{*} 60.49^{*} 56.56^{*}

TABLE S.X. Ratio between the cation and anion LJ size and energy parameters for 11 alkali halide salts.

salt	LiCl		LiI		NaCl		NaI	
Model	σ_+/σ	$\varepsilon_+/\varepsilon$	σ_+/σ	$\varepsilon_+/\varepsilon$	σ_+/σ	$\varepsilon_+/\varepsilon$	σ_+/σ	$\varepsilon_+/\varepsilon$
HMN-S	0.65	0.0015	0.54	0.0038	0.87	0.0015	0.71	0.0038
Fyta					0.59	1.0000	0.44	27.9067
JJ	0.71	0.0007	0.60	0.0007	1.01	0.0007	0.85	0.0007
RH-L	0.66	0.0050	0.56	0.0036	0.79	0.0263	0.67	0.0189
JC-T	0.29	8.9172	0.27	2.4932	0.44	14.4439	0.42	4.0385
KBFF	0.41	1.4894	0.34	3.5000	0.56	0.6809	0.46	1.6000
RDVH	0.43	1.0000	0.39	1.0000	0.43	1.0000	0.40	1.0000
MP	0.37	2.3046	0.33	1.3664	0.54	3.1279	0.48	1.8545
salt	${ m KF}$		KCl		KI		CsF	
Model	σ_+/σ	$\varepsilon_+/\varepsilon$	σ_+/σ	$\varepsilon_+/\varepsilon$	σ_+/σ	$\varepsilon_+/\varepsilon$	σ_+/σ	$\varepsilon_+/\varepsilon$
HMN-S	1.32	0.0013	1.03	0.0015	0.85	0.0038	1.51	0.0013
Fyta	0.65	162.6667	0.61	5.8290	0.46	162.6667	0.84	21.6667
JJ	1.70	0.0007	1.29	0.0007	1.07	0.0007	2.03	0.0007
RH-L	1.04	0.5015	0.84	0.2906	0.72	0.2090	1.18	1.4012
JC-T	0.63	177.4156	0.58	23.9648	0.54	6.7005	0.74	250.4011
KBFF	0.90	0.1300	0.76	0.2766	0.62	0.6500	1.12	0.0065
RDVH	0.76	1.0000	0.63	1.0000		1.0000	0.98	1.0000
MP	0.81	18.2985	0.69	4.7282	0.61	2.8033	0.87	77.5406
salt	CsCl		CsI		RbCl			
Model	σ_+/σ	$\varepsilon_+/\varepsilon$	σ_+/σ	$\varepsilon_+/\varepsilon$	σ_+/σ	$\varepsilon_+/\varepsilon$		
HMN-S	1.18	0.0015	0.97	0.0038				
Fyta	0.79	0.7764	0.60	21.6667				
JJ	1.54	0.0007	1.29	0.0007	1.39	0.0007		
RH-L	0.96	0.8119	0.82	0.5840	0.88	0.5040		
JC-T	0.68	33.8234	0.64	9.4569	0.62	37.1435		
KBFF	0.94	0.0138	0.77	0.0325	0.82	0.3191		
RDVH	0.81	1.0000	0.75	1.0000	0.74	1.0000		
MP	0.75	20.0360	0.66	11.8791	0.72	9.6123		

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