

# A Combined Wetting and Scattering Study of the Near Surface Ordering in Sugar Surfactant Based Bicontinuous Microemulsions at Hydrophilic and Hydrophobic Surfaces

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Wellert S, Stehle R, Micciulla S, Dahl M, Steitz R, Hellweg T and Holderer O (2022) A Combined Wetting and Scattering Study of the Near Surface Ordering in Sugar Surfactant Based Bicontinuous Microemulsions at Hydrophilic and Hydrophobic Surfaces. Front. Soft. Matter 2:887610. doi: 10.3389/frsfm.2022.887610 The commercial availability of natural surfactants, e.g., alkyl-oligoglucosides and the solubilization of plant and food grade oils extends the field of applications for microemulsions. To study potential effects of the confinement on the structure and dynamics inside a microemulsion, neutron reflectometry and neutron spin echo spectroscopy under grazing incidence have been used. Measurements of the contact between a bicontinuous microemulsion and a hydrophilic and hydrophobic surface are compared and show a similar wetting behavior and near surface structuring.

Keywords: microemulsion, bending elasticity, neutron scattering, neutron spin echo, neutron reflectomery

## **1 INTRODUCTION**

Several applications, such as decontamination and cleaning of sensitive surfaces Mulligan et al. (2001); Giorgi et al. (2010); Vargas-Ruiz et al. (2016), benefit from the use of mild and environmentally compatible microemulsions made from sustainable resources. The contact between a microemulsion and a solid surface is characterized by the wetting properties and the interaction potential at the surface.

This presence of a confining surface results either in surface melting or ordering of the bulk structure of the fluid in the vicinity of the surface. This is a general phenomenon previously observed for a variety of systems like metals, molecular crystals, and also colloidal systems like thermotropic liquid crystals Lang et al. (2000), dispersions Klapp et al. (2008), micellar solutions, Gerstenberg et al. (1998), Gerstenberg et al. (2002); Gerstenberg and Pedersen (2001) and microemulsions Lang (2004). To investigate the inner structure of the adhering fluid, surface sensitive scattering (GISAXS, GISANS, neutron, and X-ray reflectometry), atomic or surface force measurements Fragneto-Cusani (2001); Majewski et al. (2000); Müller-Buschbaum (2013); Petrov et al. (1997) and spectroscopic techniques are frequently used.

In case of surfactant containing fluids previous work investigated the near surface structure of sponge phases Bowers et al. (2004); Hamilton et al. (2002) and bicontinous microemulsions on planar surfaces using scattering and numerical methods Zhou et al. (1992); Lee et al. (1995); Kerscher et al. (2011).

Recently, a combination of grazing incidence geometry with neutron spin echo spectroscopy (GINSES) was established. New experimental routes for the investigation of the dynamics near solid surfaces in the time range between a few ns up to 100 ns are opened. Experiments on shorter time and length scales compared to light scattering are possible and multiple scattering can be avoided. The GINSES technique was used to study the dynamics of  $C_iE_j$  surfactant based bicontinuous microemulsions near solid surfaces as a function of the neutron penetration depth Frielinghaus et al. (2012). Furthermore, the relaxation of thermal fluctuations near solid interfaces was also investigated for other soft matter systems. In addition to phospholipid membranes Jaksch et al. (2019), these also include adsorbed microgels, e.g., Witte et al. (2021).

In the present study, the influence of the chemical nature of the confining surface on the near surface structure and dynamics of a sugar surfactant based microemulsion is addressed. In this quaternary phase system cyclohexane is used as oil phase. Cyclohexane is frequently used as solvent in chemical and industrial applications. Moreover, the majority of experimental work on microemulsions uses nonpolar or polar linear hydrocarbons while studies using nonlinear or aromatic oils in such phase systems are rare Burauer et al. (2000). Previously, the phase behavior of this quaternary system and its bicontinuous bulk structure had been characterized Wellert et al. (2011).

To the best of our knowledge, most experiments addressing the interactions between solid surfaces and microemulsions were carried out either on non-ionic C<sub>i</sub>E<sub>i</sub>-type or ionic surfactants. Microemulsions coexisting with excess phases in the SDHStoluene-saline water system were measured using contact angle measurements to study the solid-liquid-liquid wetting behavior Stammitti-Scarpone and Acosta (2019). Only few results from microemulsions made of bio-inspired surfactants were published. For several reasons, such experiments are crucial. Beside the differences in the composition of the amphiphilic interface and its tuning mechanism also the adsorption properties of these surfactants differ significantly Matsubara et al. (2009); Grosse and Estel (2000). For example, sugar surfactants as n-dodecyl- $\beta$ -d maltoside adsorb well on hydrophilic surfaces as alumina, hematite, and titania but adsorb less on silica. Ethoxylated surfactants show the opposite behavior (Zhang et al. (2002); Lu et al. (2007). Here, we address the characterization of wetting behavior, inner structure and dynamics near a solid substrate of sugar-surfactant based microemulsions in the vicinity of hydrophilic and hydrophobic surfaces. For this purpose, we combine contact angle measurements and the surface sensitivity of neutron reflectometry and GINSES.

## 2 MATERIALS

As the amphiphilic components of the microemulsions under investigation Glucopon 220 and the medium chain alcohol pentanol (Sigma-Aldrich) were used. was used. of alkylpolyglucoside  $C_{8/10}G_{1.3}$  (Cognis, Germany) Glucopon 220 is a technical grade aqueous stock solution (43%, aq.) of the alkylpolyglucoside  $C_{8/10}G_{1.3}$  (Cognis, Germany). It was freeze dried to a residual water content less than 1%, determined by Karl-Fischer titration.

The bulk phases of the microemulsions were formed by cyclohexane (Sigma-Aldrich) and  $D_2O$  (purity  $\geq$ 98%, isotopic purity  $\geq$ 99.9%, Euriso-Top, France).

For cleaning the silicon substrates with RCA solution  $NH_4OH$ at 30%, aq. and  $H_2O_2$  (Sigma-Aldrich) were used. Water was puriefied with a Millipore Milli-Q system. For the surface scattering experiments, silicon blocks (5 × 8 × 1.5 cm) were purchased from Siliciumbearbeitung Andrea Holm, Tann, Germany. For the functionalization of the silicon substrates, dichlorodimethylsilane (HMDS), anhydrous chloroform, and methanol purchased from Sigma-Aldrich were used.

# **3 EXPERIMENTAL SECTION**

## **3.1 Bicontinuous Microemulsion**

The bicontinuous microemulsions investigated in this work were prepared according to the composition given by the quaternary phase diagram shown in **Supplementary Figure S1** in the supporting material. At a constant cyclohexane to water ratio of  $\alpha = 0.5$  the single phase domain is depicted as a function of the  $C_{8/10}G_{1.3}$  (Glucopon 220) weight fraction  $\gamma$  and the total pentanol weight fraction  $\delta$ . A more detailled discussion of the phase behavior can be found in the literature Wellert et al. (2011). At a surfactant concentration of  $\gamma = 0.124$  the bicontinuous region appears. With increasing surfactant content above  $\gamma = 0.25$  the one phase region passes into the lamellar phase. For the experiments a bicontinuous sample at  $\gamma = 0.21$ , safely far from the phase boundaries, was chosen.

# **3.2 Surface Preparation and Functionalization**

The silicon blocks were cleaned with RCA solution (5:1:1 of water,  $NH_4OH$  at 29%,  $H_2O_2$  at 30%) at 72 °C for 10 min, then rinsed several times with water and dried in a nitrogen stream.

For functionalization the silicon blocks or wafers were transferred from water to chloroform by subsequent ultrasonication in mixtures of CHCl<sub>3</sub>/CH<sub>3</sub>OH = 1:1, 3:1 and pure CHCl<sub>3</sub> for 10 min. Afterwards, the silicon substrates were put in an upright position in a desiccator with a residuum of approximately 6  $\mu$ L) of HMDS in a vacuum desiccator at ambient temperature for 24 h. Then the substrates were transferred from chloroform to methanol by subsequent ultrasonication for 5 min in solvent mixtures of CHCl<sub>3</sub>:CH<sub>3</sub>OH from 1:0  $\rightarrow$  0:1. The substrates were rinsed several times with water and dried in a nitrogen stream.

RCA 1 cleaned silicon blocks (static water contact angle  $\leq 10^{\circ}$ ), were used as hydrophilic surface, whereas HMDS-modified silicon blocks (static water contact angle ( $80 \pm 10^{\circ}$ ) were used as hydrophobic surface. The RCA cleaning protocol was applied just before sealing the substrate inside the solid/liquid cell dedicated to NR and GINSES measurements. In the remainder of this paper, hydrophilic and hydrophobic surfaces are termed h-Si respectively hp-Si.

#### **3.3 Surface Tension Measurements**

The surface tension of the pure MilliQ-type water, cyclohexane, aqueous sugar surfactant solution (c ( $C_{8/10}G_{1.3}$ ) in H<sub>2</sub>O of 20 mmol/ml) and the microemulsion were determined with a force tensiometer DCAT 11 (Dataphysics, Germany) applying the du Noüy ring method. Prior to each measurement, the platinum ring was flamed to remove any organic residuals. For each sample, 10 measurements were carried out and averaged at 25°C.

#### 3.4 Contact Angle Measurements

Contact angles were measured with a drop contour analysis of sessile drops on the planar hydrophilic (h-Si) and hydrophobic (hp-Si) substrates using the OCA20 contact angle goniometer (Dataphysics, Germany). Prior to the measurements, the sample surfaces and the sessile droplets were placed in a quartz sample cell in a saturated vapor phase using the investigated liquids. Measurements of the surface energy components of h-Si and hp-Si substrates are shown in **Supplementary Figure S2** and described in the supporting information.

#### 3.5 Neutron Reflectometry

Neutron reflectivity was measured at the neutron reflectometer V6 located at the NL4 neutron guide of the medium flux research reactor at the Helmholtz Zentrum Berlin (Germany). The neutron wavelength of 4.66 Å was selected by a graphite monochromator and filtered by a liquid nitrogen cooled Be filter. The beam profile was  $1 \times 40$  mm which sets the instrument resolution to  $\Delta Q = 2 \times 10^{-3} \text{ Å}^{-1}$ . The sample cell consisted of a Teflon trough closed with the single crystal silicon block serving as the solid interface in a thermostated aluminium housing. This cell was placed horizontally in the collimated neutron beam. Both front ends of the cell were covered with cadmium shielding to prevent bulk phase scattering.

Reflectivity curves were measured in  $\theta/2\theta$ -geometry and recorded using a position sensitive detector. The scattering length densitites (SLDs) of the microemulsion components are  $6.38 \times 10^{-6} \text{ Å}^{-2}$ ,  $6.72 \times 10^{-6} \text{ Å}^{-2}$ ,  $0.34 \times 10^{-6} \text{ Å}^{-2}$ , and  $-0.32 \times 10^{-6} \text{ Å}^{-2}$  for D<sub>2</sub>O, deuterated cyclohexane (C<sub>6</sub>D<sub>12</sub>), C<sub>8/10</sub>G<sub>1.3</sub> and pentanol (C<sub>5</sub>H<sub>11</sub>OH). Additionally, water with a SLD of  $-0.53 \times 10^{-6} \text{ Å}^{-2}$  was used in a D<sub>2</sub>O/H<sub>2</sub>O mixture to adjust the bulk phase SLD. To obtain reliable fit results, the critical edge of total reflection of a reflectivity curve has to be within the probed Q<sub>z</sub> range. To achieve this, the contrast between oil phase and water phase was accordingly tuned, while still maintaining a SLD step towards higher SLD from Si (2.1 × 10<sup>-6</sup> Å<sup>-2</sup>) to the average SLD of the microemulsion. Raw data were normalized and footprint corrected with the software provided by the HZB.

The measured reflectivity curves were analyzed with the fitting routine Parratt 32. This approach determines the neutron reflectivity from planar surfaces using a calculation based on Parratt's recursion scheme for stratified media. An oscillating SLD profile of the form

$$\rho(z) = Ae^{\left(-z/\xi_z\right)} \cos\left(\left(\frac{2\pi z}{D_z}\right) + \phi\right) \tag{1}$$

was used in this approach. The SLD model assumes alternating layers of oil and water with a domain size  $D_z$  perpendicular to the confining substrate. The ordered structure vanishes according to a correlation length  $\xi_z$ . Based on Ginzburg–Landau theory this order parameter profile was obtained by including additional surface fields in the free energy functional Lee et al. (1995); Gompper and Schick (1990); Chernov and Mikheev (1988). Similar theoretical concepts are used for the description of planar amphiphilic membranes adhered to a planar solid substrate Charitat et al. (2008).

## 3.6 Neutron Spin Echo Spectroscopy in Transmission (NSE) and Reflection Geometry Under Grazing Incidence (GINSES)

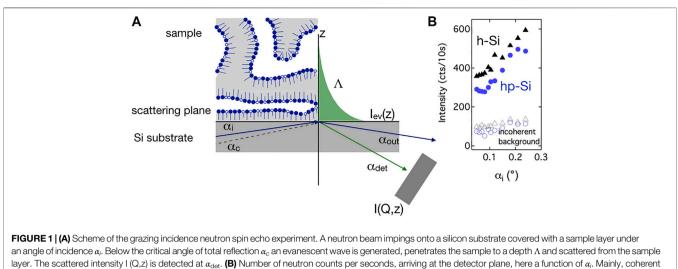
All NSE measurements were carried out on the J-NSE instrument at the FRM II neutron source in Garching, Germany Holderer et al. (2008). A neutron wavelength of 8 Å was used. All samples were measured at  $25^{\circ}$  C.

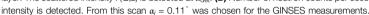
The preferred contrast for NSE is the "film contrast", where only the interface layer contains hydrogenated materials and oil and water phase are deuterated with as little hydrogen as possible. Interface fluctuations can be best observed in this contrast. The measurements on the bulk sample were done in transmission geometry using standard quartz sample cells. The neutron path length in the Hellma quartz cells was 2 mm. The bulk sample was measured at  $Q = 0.08 \text{ Å}^{-1}$  for reference.

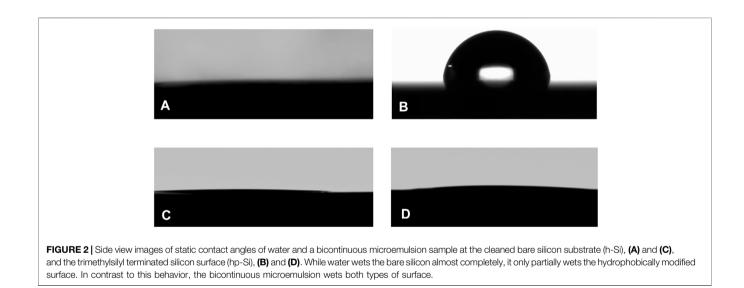
Neutron spin echo spectroscopy under grazing incidence requires a well defined angle of incidence  $\alpha_i$  to provide a defined depth of neutron penetration  $\Lambda$  into the sample Frielinghaus et al. (2012); Holderer et al. (2014). Therefore the entrance aperture is reduced to 2 mm. **Figure 1** illustrates the GINSES scattering geometry. The neutron beam enters the sample from the silicon block and is reflected at the interface. The critical angle of total reflection  $\alpha_c = \lambda \sqrt{\Delta \rho / \pi}$  depends on the chosen neutron wavelength  $\lambda$  and the contrast  $\Delta \rho$  between the silicon block and the average SLD of the microemulsion. At an angle of incidence  $\alpha_i \leq \alpha_c$  an evanescent neutron wave of intensity  $I_{ev}(z)$  exponentially decays into the sample over a characteristic depth  $\Lambda$  with

$$\Lambda = \frac{1}{\sqrt{4\pi\Delta\rho \left(1 - \alpha_i^2 / \alpha_c^2\right)}}.$$
 (2)

Plot b) in **Figure 1** shows the spin coherent and incoherent scattered intensities at  $Q = 0.08 \text{ Å}^{-1}$  as a function of  $\alpha_i$  for both surfaces. In the spin-up configuration (all polarization flippers off) no spin manipulation occurs and mainly the coherent contribution is measured. Contrary, the spin-down configuration (180°-polarization flip at the sample position) is dominated by incoherent scattering. At increasing angle of incidence  $\alpha_i$  the scattered intensities slightly increase until  $\alpha_c$  is reached. From the SLDs of deuterated cyclohexane and D<sub>2</sub>O, constituting the bulk phases of the bicontinuous microemulsion, a value of  $\alpha_c \approx 0.15^\circ$  was estimated. However, already below  $\alpha_c$  the coherent scattered intensities from the h-Si and hp-Si interfaces







increase steeply. An angle of incidence  $\alpha_i = 0.11^{\circ}$  close to the onset of this increase was chosen in the grazing incidence measurements. For this combination of  $\alpha_c$  and chosen  $\alpha_i$ , a neutron penetration  $\Lambda \approx 70$  nm was estimated using **Eq. 2**. This corresponds to a penetration into the sample over a distance of 4–5 domains. The scattered intensity was measured at Q = 0.08 Å<sup>-1</sup>.

## **4 RESULTS AND DISCUSSION**

### 4.1 Wetting Behavior

To study the influence of the chemical composition of a confining substrate on the near surface dynamics of bicontinuous microemulsions, freshly cleaned silicon surfaces (h-Si) and HMDS-modified silicon surfaces (hp-Si)

were used as hydrophilic and hydrophobic surfaces, respectively. For h-Si the measured surface energy was  $\gamma^{\text{SE}}=(69 \pm 3) \text{ mN/m}$  with  $\gamma_p^{\text{SE}}=(48 \pm 2) \text{ mN/m}$  and  $\gamma_d^{\text{SE}}=(20 \pm 2) \text{ mN/m}$  being the polar and unpolar components of the surface energy. In the case of hp-Si, the surface energy was  $\gamma^{\text{SE}}=(30 \pm 1) \text{ mN/m}$  with  $\gamma_p^{\text{SE}}=(8 \pm 1) \text{ mN/m}$  and  $\gamma_d^{\text{SE}}=(22 \pm 1) \text{ mN/m}$ . Details of this measurements can be found in the supporting information in **Supplementary Figure S3**. The vapor deposition chemically changes the hydroxylated silicon surface to a trimethylsilyl terminated surface with  $((\text{CH}_3)_3\text{Si-O})$  groups at the outermost layer. This is reflected in the drastic change of the wetting behavior of water. Since water spreads over the freshly cleaned silicon surface without a detectable contact angle, at the trimethylsilyl terminated silicon surface values between 70° an 90° were observed, depending on the coverage of the surface with

**TABLE 1** | Surface tensions  $\sigma$  at 25°C and contact angles of the bulk fluid components and the microemulsion at the hydrophilic (h-Si) and hydrophobic (hp-Si) silicon substrates.

Compound	σ (mN/m)	CA <sub>h-Si</sub> (°)	CA <sub>hp-Si</sub> (°)
Water	71.4 ± 0.1	≈0	70–90
Cyclohexane	24.2 ± 0.1	≈0	≈0
Glucopon 220 solution	28.2 ± 0.1	≈0	26 ± 2
Microemulsion	$24.6 \pm 0.2$	≈0	≤7

trimethylsilyl groups. In case of the bicontinous microemulsion, wetting with a very low contact angle below the detection limit of the optical contact angle goniometer was observed. This is illustrated by the images in Figure 2. Additionally, an aqueous solution of Glucopon 220 with a concentration of approximately 20 mmol/L was studied. On hp-Si a contact angle of  $(26 \pm 2)^{\circ}$  was detected while h-Si was fully wetted by the surfactant solution. The results for the static contact angles were compiled in Table 1. The static contact angle data are shown in Supplementary Figure S2. In addition to the static contact angles, also the measured surface tensions of pure water, the surfactant solution, cyclohexane and the bicontinuous microemulsion are shown in Table 1. These values suggests that due to its low surface tension, a significant amount of cyclohexane is present at the airliquid interface of the microemulsion which is the reason for the same wetting behavior of the microemulsion on both surfaces, hp-Si and h-Si.

## 4.2 Near Surface Structure of the Bicontinuous Microemulsion at the Hydrophilic and Hydrophobic Surface

The bicontinuous microemulsions show the same wetting behavior at the hp-Si and h-Si surfaces, as discussed in Section 4.1. This raises the question of whether there are differences in the microemulsion structure in the immediate proximity of the solid-liquid interfaces. Figure 3 shows the reflectivity curves of the investigated bicontinuous microemulsion in contact with the hydrophilic and hydrophobic surfaces. The critical edge of total reflection was shifted into the  $Q_z$  range accessible for the instrument mixing  $H_2O/D_2O$  to tune the SLD of the aqueous phase. The reflectivity curves show a broad but distinct Bragg peak which is slightly more pronounced and shifted to lower  $Q_z$  in case of hp-Si. The Bragg peak originates from the structuring of the microemulsion at the solid/liquid interface. The reflectivity data can be described by using a oscillating variation of the SLD related to alternating oil and water rich regions of the bicontinuous structure given by Eq. 1.

The inset in **Figure 3** shows the SLD profiles for the bicontinuous microemulsion on both, h-Si and hp-Si substrates. Starting from negative z, the profile corresponds to the silicon substrate, followed by the silicon oxide (thickness of 15 Å) and, in the case of the hp-Si, a negative SLD value for the trimethylsilyl surface layer. A similar oscillation behavior was

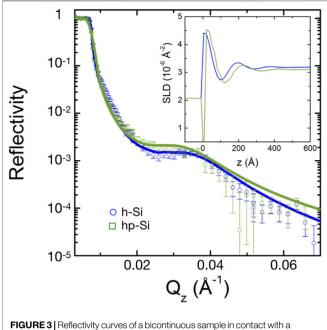


FIGURE 3 Reflectivity curves of a bicontinuous sample in contact with a hydrophilic and hydrophobized silicon surface. The lines are fits to the data calculated from the SLD profile given by 1 and shown in the inset.

observable for both measurements. The SLD oscillations decay over two oscillations to the bulk value at  $z \to \infty$ .

From the interpretation of these results, one can say, that the bicontinous bulk structure remains unchanged up to the immediate vicinity of the h-Si and hp-Si substrate. For the layer in direct contact with the substrate, a very similar SLD profile was obtained, which suggests that the structure of the microemulsion in the vicinity of either a hydrophilic or a hydrophobic surface is governed by the same ordering of the bulk phases. This corresponds well to the results of the contact angle measurements.

However, since these measurements determine the reflectivity without additional information of the phase of the complex reflection coefficient, in the kinematical approxiamation the near surface structure could only be determined to a remaining ambiguity. Hence, no reliable conclusion can be drawn about the existence and composition of an adsorbed surfactant layer at the surface. The "phase problem" in neutron reflectometry, i.e., the loss of phase information due to the measurement of reflected intensities, could be addressed in principal by measuring the same sample in different contrasts by isotope substitution or by introducing a known reference layer at the interface and varying the interfacial contrast by polarization analysis.

However, the use of the oscillating SLD profile given by Eq. 1 enables the estimation of the domain size  $D_z$  and the correlation length  $\xi_z$  perpendicular to the confining substrate. These results and the corresponding values of the bulk microemulsion,  $d_{TS}$  and  $\xi_{TS}$ , measured with SANS are compiled in **Table 2**. The sizes of

TABLE 2 Summary of domain sizes D, correlation lengths  $\xi$  determined from SANS in bulk and neutron reflectometry at the solid-liquid interface. The renormalized and bare bending elasticities KSANS and KDARE obtained and calculated from the fit results of the reflectometry and previous SANS measurements. Data of the SANS measurements are taken from reference (Wellert et al. (2011)).

Parameter	Bulk Phase (SANS)	h-Si surface	hp-Si surface
Correlation length $\xi$ (Å)	70 ± 1	111 ± 6	90 ± 5
Renormalized bending elasticity $\kappa_{SANS}$ (k <sub>B</sub> T)	$0.38 \pm 0.02$	$0.54 \pm 0.08$	0.42 ± 0.07
Bare bending elasticity $\kappa_{bare}$ (k <sub>B</sub> T)	$0.82 \pm 0.10$	$0.90 \pm 0.22$	$0.90 \pm 0.22$

the oil and water domains from the measurements at the h-Si and the hp-Si surfaces are similar. Compared to the bulk phase the correlation lengths are larger in case of the microemulsions in contact with the h-Si and the hp-Si surfaces. This difference suggests, that a small effect of the planar interface on the inner structure of the microemulsion exists. The SLD profiles of both measurements have a very similar course. This suggests a similar composition of the microemulsions in the near surface region of the two solid liquid interfaces.

The differences between the results from the bulk measurements and those at the interface must be interpreted with caution. Due to the pronounced maximum in the scattering curve in bulk, the peak position can be found more accurately in the fit than in the analysis of the reflectivity data where the maximum is strongly broadened. It can be deduced from the results that the domain sizes are modified, if at all, only to a small extent by the presence of the interfaces.

In a SANS measurement, the renormalized and the bare bending elasticity constants  $\kappa_{SANS}$  and  $\kappa_{bare}$  of the surfactant membrane in the microemulsion Gompper et al. (2001); Gompper and Kroll (1998),

10

and

$$\frac{\kappa_{SANS}}{k_B T} = \frac{10\sqrt{3}\pi}{64} \frac{\xi_{TS}}{d_{TS}}$$
(3)

$$\frac{\kappa_{bare}}{k_B T} = \frac{\kappa_{SANS}}{k_B T} + \frac{3}{4\pi} \ln\left(\frac{d_{TS}}{2l_c}\right) \tag{4}$$

can be calculated from the structural parameters (domain size  $d_{TS}$  and correlation length  $\xi_{TS}$ ) obtained from the Teubner–Strey analysis of the SANS data.  $l_c$  represents the thickness of the surfactant membrane. The same approach was applied to estimate the values of the renormalized and bare bending elasticity constants from the neutron reflectometry data. From the domain size  $D_z$  and correlation length  $\xi_z$  from reflectometry, the bendig rigidity has been calculated with Eq. 3. In particular the comparison between the hydrophilic and hydrophobic interface is possible with this approach, and it also gives a strong hint that the membrane elasticity is not significantly altered by the interface layer, as mentioned above.

The renormalized bending modulus  $\kappa_{SANS}$  and the bare bending modulus  $\kappa_{bare}$  characterize the elastic properties of the amphiphilic film at different length scales. On the length scale of the oil and water domains,  $\kappa_{SANS}$  is the effective bending modulus

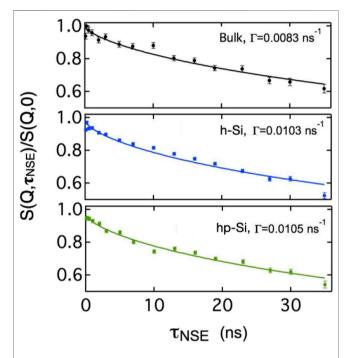


FIGURE 4 | Normalized intermediate scattering functions of a bulk contrast microemulsion measured in top: bulk in transmission geometry, middle: in contact with a bare silicon substrate under grazing incidence and bottom: in contact with a hydrophobized silicon substrate under grazing incidence. The solid lines are fits to a stretched single exponential function with stretching exponent 2/3 as suggested by the Zilman-Granek approach for undulating membrane patches (The relative errors of the resulting relaxation rates are 9%, 5%, and 7%,)

of the membrane. The bending motion of the membrane is superimposed with thermal fluctuations at smaller scales. Strong thermal fluctuations result in a larger deformation of the membrane Gompper et al. (2001); Pieruschka et al. (1995). The effective thickness of the surfactant membrane between the oil and water domains  $l_c$  determines the bare bending elasticity  $\kappa_{bare}$  Gompper and Kroll (1998); Morse (1994). The values of  $\kappa_{SANS}$  are rather small, between  $0.54 k_B T$  and  $0.42 k_B T$  and support the conclusion of a highly flexible and very soft interface. These values are smaller compared to the results from previous experiments with bicontinuous microemulsions from the ternary phase system decane/ $H_2O/C_{10}E_4$  at hydrophilic surfaces. For these microemulsions, the formation of a lamellar

ordering in the vicinity of the surface was found. The spatial expansion of this ordering away from the substrate extended up to a few domain sizes.

#### 4.3 Bulk and Near Surface Dynamics

The membrane dynamics as measured with NSE spectroscopy can be described by calculating the intermediate scattering function S (Q, $\tau_{NSE}$ ) of fluctuating membrane patches Zilman and Granek (1996); Mihailescu et al. (2001). Before applying often used approximations for large bending rigidities, the intermediate scattering function is:

$$S(\vec{Q}, \tau_{NSE}) \propto \left\langle \int d^2r \int d^2r' \exp(i\vec{Q}_{xy}(\vec{r} - \vec{r}')) + \exp\left(-\frac{k_B T}{4\pi^2 \kappa_{bare}} Q_{z}^2 \int_{k_{min}}^{k_{max}} \frac{dk}{k^4} \left[1 - e^{i\vec{k} \cdot (\vec{r} - \vec{r}') - \omega(k)\tau}\right] \right) \right\rangle$$
(5)

This can be approximated by a stretched exponential function which shows typically a rate  $\Gamma \propto Q^3$  and a stretching exponent  $\beta$  and yields by this information important hints if membrane dynamics is observed and what the direct effects of the interfacial confinement on the membrane dynamics are.

In **Figure 4** three plots of normalized intermediate scattering functions (ISF) are shown. The ISF at the top was measured from the bulk sample in transmission geometry measured at a momentum transfer  $Q = 0.08 \text{ Å}^{-1}$ . Below this, two ISF's were measured under grazing incidence at  $Q_{GINSES} \approx Q_z = 0.08 \text{ Å}^{-1}$ . The data from the bicontinuous microemulsion in contact with the hydrophilic (h-Si) sample are given in the middle graph. Plotted in the bottom graph is the ISF of the microemulsion adhered to the trimethylsilyl terminated surface (hp-Si). All three curves show a continuous, smooth decay within the Fourier time range  $\tau_{NSE}$ .

As a first step in the analysis, a stretched exponential with a stretching exponent  $\beta = 2/3$ ,

$$S(Q, \tau_{NSE}) / S(Q, 0) = Aexp(-(\Gamma \tau_{NSE})^{\beta}).$$
(6)

was fitted to the data. The decay rate  $\Gamma$  measures the relaxation of thermally excited fluctuations of the surfactant film. The resulting relaxation rates  $\Gamma$  are given in the plots. A comparison shows, that the relaxations close to the surface are slightly faster than in bulk but are the same for both types of surface within the precision of the fit.

A detailed analysis was achieved by applying the full integral evaluation, required for small bending rigidities of the order of 1 k<sub>B</sub>T, usually observed for bicontinuous microemulsions. In this model, the bending rigidity  $\kappa$  is the only fitting parameter.

Bulk microemulsions have a dispersion relation  $\omega(k) = \frac{\kappa_{borr}k^3}{4\eta}$ with the undulation wave number k and the solvent viscosity  $\eta$ . For near surface dynamics, long wavelength undulations show a different dispersion relation  $\omega(k) \propto k^2$  due to interactions with the rigid wall Frielinghaus et al. (2012) Seifert (1994); Kraus and Seifert (1994). The interaction potential between the microemulsion and the confining wall introduces another length scale, which is related to  $\bar{l}$ , the distance between the first membrane and the wall. GINSES experiments allow to fit this length if using the full integral version of S (Q,  $\tau_{NSE}$ ). This has been measured with GINSES for bicontinuous microemulsions Frielinghaus et al. (2012); Lipfert et al. (2014) and requires a fit with the full Zilman-Granek theory involving the numerical integration of all undulation waves of the membrane patches.

First, the bulk microemulsion has been fitted to determine the required parameters of the fitting function, especially the integration limit for long wavelength undulations, which is four times the correlation length determined from SANS. The bending rigidity has been fixed to  $\kappa_{bare} = 0.82 k_B T$  to reduce the number of fitting parameters. This corresponds to the renormalized value of  $\kappa_{SANS} = \kappa_{bare} - \frac{\alpha}{4\pi} \ln(d_{TS}/2l_c)$  with the molecular length  $l_c$  of about 12 Å and the distance  $d_{TS}$  between domains in the microemulsion and  $\alpha \approx 3$ . Additional corrections would be necessary from the influence of the saddle splay modulus  $\bar{\kappa}$  which can be neglected for the purpose of fixing the fit parameters for the interface layer Holderer et al. (2013).

In the next step, the fit of the microemulsion dynamics at both types of interfaces, the parameters were fixed by the bulk microemulsion except the characteristic distance  $\bar{l}$  between first layer and wall. The integration limit for the long wavelength undulations determined in the bulk by  $\xi$  is set to infinity. In this case,  $\bar{l}$ , measuring the wall-membrane interaction remains as the only fitting parameter to describe the difference between bulk and interface dynamics. Details can be found in literature Frielinghaus et al. (2012). The fit yields  $\bar{l} = (24 \pm 4)$  Å and  $(33 \pm 4)$  Å for the h-Si and hp-Si interface, respectively. Within the precision of the experiment, both values are the same. The fits to the data are shown in **Supplementary Figure S4** in the supporting information. Additionally, the fit parameters are summarized in **Supplementary Figure S5**.

These results agree with the findings from the neutron reflectometry measurements and suggest minor differences in the coupling of the investigated bicontinuous microemulsion to both surfaces. Probably, the microemulsion bulk structure adheres to both solid interfaces without a distinct extended lamellar transition region. In a previous experiment with ternary bicontinuous microemulsions in the decane/H<sub>2</sub>O/C<sub>10</sub>E<sub>4</sub> phase system in contact with a hydrophilic silicon substrate (h-Si) faster relaxation was found in the vicinity of the interface. On the other hand, for the contact between clay particle surfaces and bicontinuous microemulsion of the same phase system no difference between bulk and near surface relaxations in the microemulsions were found, which also was attributed to the absence of a lamellar transition region Frielinghaus et al. (2012).

The results of these GINSES experiments indicate a significant influence of the interaction between surfactant molecules and confining surface on the formation of an adjacent lamellar structure. Comparative studies on the adsorption of nonionic ethoxylated surfactants and of sugar surfactants on a variety of substrates revealed significant differences. The adsorption of a surfacatant is influenced by the chemical composition of the adsorbent. Moreover, although in both cases hydrogen bonding mediates the adhesion of the surfactant, at silicate surfaces ethoxylated surfactants were found to adsorb stronger than sugar surfactants Zhang et al. (2006); Zhang and Somasundaran (2006) which might support the formation of lamellar ordering in a  $C_iE_j$  based bicontinuous microemulsion.

## **5 CONCLUSION**

Wetting and neutron scattering measurements were combined to investigate the influence of the chemical composition of a solid surface on the interfacial structure and dynamics of a sugar surfactant-based bicontinuous microemulsion. Due to their environmental compatibility and temperature-independent phase behavior, sugar surfactant-based microemulsions are suitable for applications in temperature-variable environments, for example.

Bicontinuous microemulsions with cyclohexane as the nonpolar bulk phase were found to completely wet both, hydrophilic and hydrophobic surfaces. Investigation of the interfacial structure with neutron reflectometry showed nearly identical structuring at both interfaces in terms of similar sizes of the oil and water domains in the bulk and at the planar surfaces and a similar SLD profile indicating a similar composition of the layers in close proximity to the solid liquid interface. Slightly larger correlation lengths were found at the planar substrates.

Similar results were observed for the interfacial dynamics in these microemulsions. Here, deviations from the relaxation of thermal fluctuations in the bulk phase occur. An influence of the hydrophobicity of the surface or the chemical structure was not observed. In combination with the weak adsorption of the sugar surfactants, the formation of a extended lamellar structure close to the interface is suppressed or strongly hindered.

For technical applications where good wettability of hydrophobic and hydrophilic surfaces is required, such sugar surfactant-based microemulsions thus offer a promising medium, for example for extraction applications.

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## DATA AVAILABILITY STATEMENT

The raw data supporting the conclusion of this article will be made available by the authors, without undue reservation.

### AUTHOR CONTRIBUTIONS

Conception and idea from SW, TH, and ROS. Sample preparation and contact angle measurements SM, MD, SW, SW, RAS, and ROS conducted and analyzed the neutron reflectometry experiments. SW, SM, and OH conducted and analyzed the neutron spin echo experiments. Further data evaluation and writing of the paper by SW and OH.

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### SUPPLEMENTARY MATERIAL

The Supplementary Material for this article can be found online at: https://www.frontiersin.org/articles/10.3389/frsfm.2022.887610/full#supplementary-material

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