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Converting high dimensional complex networks to lower dimensional ones preserving synchronization features

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Abstract –Studying the stability of synchronization of coupled oscillators is one of the prominent topics in network science. However, in most cases, the computational cost of complex network analysis is challenging because they consist of a large number of nodes. This study includes overcoming this obstacle by presenting a method for reducing the dimension of a large-scale network, while keeping the complete region of stable synchronization unchanged. To this aim, the first and last non-zero eigenvalues of the Laplacian matrix of a large network are preserved using the eigen-decomposition method and, Gram-Schmidt orthogonalization. The method is only applicable to undirected networks and the result is a weighted undirected network with smaller size. The reduction method is studied in a large-scale a small-world network of Sprott-B oscillators. The results show that the trend of the synchronization error is well maintained after node reduction for different coupling schemes.

Introduction. – The study of structural and dynamical features of real-world networks is facilitated using com-2 plex networks in different fields such as biology [1,2], neuroscience [3,4], ecology [5,6], and social science [7,8]. Synchronization is an important topic in complex networks [9–11]. Different types of synchronization have been found 6 there, including complete [12], phase [13], cluster [14–16], 7 explosive [17], and lag synchronization [18]. These syn-8 chronized states can emerge as the effect of the static or time-varying interactions in either attractive or repulsive 10 couplings [19]. Moreover, enormous effort has been de-11 voted to the controllability and observability of synchro-12 nized complex networks [20], improving the synchroniz-13 ability [21, 22] and robustness of synchronization [23]. 14

¹⁵ Most real-world systems can be better modeled by com-¹⁶ plex networks or even wih considering higher-order inter-¹⁷ actions [24]. These models may contain many nodes, mak-¹⁸ ing their analysis difficult and costly. Therefore, any effi-¹⁹ cient reduction of the size is of interest. One of the basic ²⁰ methods to decrease network size is graph partitioning. Various criteria exist whose persistence has been consid-21 ered in reducing the network nodes. For instance, authors 22 in [25] seek to keep some physical properties of the net-23 work after node reduction. Graph partitioning methods 24 are mostly considered as non-deterministic polynomial-25 time problems [26], which cannot be solved in polynomial 26 time. Therefore, researchers have tried to find other meth-27 ods to reduce the network size. For example, Bona et al. 28 [27] proposed a reduced model for the public transporta-29 tion complex network with a long sequence of 2-degree 30 nodes and some hubs. Despite removing 2-degree nodes, 31 the reduced network has the same topological characteris-32 tics and skeleton as the original one. Besides, it was shown 33 that this reduction increases the network cluster coefficient 34 and the average degree while decreasing the path length. 35

Recently, different methods such as Spectral Coarse-Graining [28] and a Search Algorithm to Dimension Reduction [29] have been proposed. These algorithms decrease the dimension of the Laplacian matrix of the graph, while preserving some specific features of the parent net-

work to keep synchronization. Whereas Spectral Coarse-41 Graining [28] iteratively reduces the dimension by merging 42 the nodes, the Search Algorithm [29] can effectively reduce 43 the number of nodes through a fast search. Another sys-44 tematic approach for size reduction has been taken into 45 account recently. In 2020, Thibeault et al. developed 46 the Dynamics Approximate Reduction Technique to sim-47 plify a complex network [30]. Their method, which was 48 based on spectral graph theory, enabled the prediction of 49 the synchronization regimes of phase oscillators in large-50 scale networks by using dominant eigenvectors features. 51 In this method, the reduced network size is not arbitrary 52 and depends on the number of the network's communi-53 ties. In [31], the authors have reduced the dimension of a 54 non-locally coupled network by projecting the network dy-55 namics onto the subspace that corresponds to the unstable 56 eigenvalues of the linear part of the network. 57

In this paper, we introduce a novel approach to reduce 58 the size of a complex undirected network with preserv-59 ing its synchronization pattern. The key point for main-60 taining the synchronization stability of a network is to 61 keep the eigenvalues of the Laplacian matrix that affect 62 the synchronization within the master stability function 63 approach. To this end, the eigen-decomposition and the 64 Gram-Schmidt methods are utilized, and a smaller adja-65 cency matrix which is weighted, is obtained. 66

The paper is organized as follows: First, the dimen-67 sion reduction method is described in Section 2 in detail. 68 Then, a large-scale network of chaotic Sprott-B systems is 69 analyzed, and the preservation of synchronization pattern 70 after reduction is checked. The results are presented in 71 Section 3. Finally, the conclusions of the paper are given 72 in Section 4. 73

Dimension reduction method. - This section de-74 scribes the method used to reduce the dimension of a large 75 undirected network to a smaller one. The aim is to pre-76 serve the synchronization pattern of the large-scale net-77 work after dimension reduction. It has been shown that 78 the stability of synchronization in networks relies on the 79 coupling topology [32]. According to the master stability 80 function method [33], the region of stable synchronization 81 depends on the eigenvalues of the connectivity matrix of 82 the graph. Here, the reduction method is based on obtain-83 ing a reduced connectivity matrix with desired eigenvalues 84 which are those involved in determining the synchroniza-85 tion stability region. The eigen-decomposition factoriza-86 tion is used for finding this reduced connectivity matrix. 87

Master stability function. The master stability function (MSF) [33] is a method for finding the local stability of synchronization. The description of this approach is given in the following.

It is supposed that N identical oscillators with the individual dynamics of F(.) are linearly coupled by the overall coupling strength d through a Laplacian connection matrix G. For the oscillator i, one can write

$$\dot{X}_i = F(X_i) - d\sum_{j=1}^N G_{ij}H(X_j), \quad i = 1, 2, \dots, N$$
 (1)

where H indicates the coupling function. When all oscillators lie in the synchronization manifold, i.e., $X_1 =$ $X_2 = \ldots = X_N = X_s$, the linearization of Eq. (1) around the synchronized solution X_s is defined as the variational equation and can be written as

$$\dot{\eta}_{l} = [DF(X_{s}) - \alpha_{l}DH(X_{s})]\eta_{l}, \qquad l = 1, 2, \dots, N$$
 (2)

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in which $\alpha_l = d\lambda_l$, where λ_l is the l^{th} eigenvalue of the matrix G. Also, DH and DF are the Jacobian matrixes of H and F, respectively. The variational equation (Eq. (2)) determines the stability of synchronization, which can be found by calculating its maximum Lyapunov exponent. The maximum Lyapunov exponent (A) of Eq. (2) as a function of $\alpha = d\lambda$ is known as the master stability function (MSF). Considering a connected and undirected network, the first eigenvalue of G is zero $(\lambda_1 = 0)$, which is along the synchronization manifold. The other eigenvalues are sorted assendingly $\lambda_2 \leq \lambda_3 \leq \ldots \leq \lambda_N$. When $\Lambda < 0$ for all eigenvalues $\lambda_i, i = 2, \ldots, N$ of the Laplacian matrix, all the nodes of the network oscillate in complete 100 synchrony. 101

Huang et al. [34] proposed a general scheme for categorizing the MSFs and introduced four classes. The classification is based on the number of zero-crossing points of the master stability function curve versus α , such that Γ_k represents a class in which $\Lambda(\alpha)$ crosses the zero k times. In case the synchronization cannot be reached for any α value, the master stability function has no zero-crossing point and is classified as Γ_0 (Fig. 1a). The master stability function with only one zero-crossing point, α_{\min} , is known as class Γ_1 , which is shown in Fig. 1b. Sorting the eigenvalues of the Laplacian matrix (G) in ascending order (i.e., $\lambda_1 = 0$), the synchronization manifold of this class is stable if

$$\alpha_{\min} < d\lambda_2 \le d\lambda_3 \le \dots \le d\lambda_N \tag{3}$$

holds. Hence, choosing the coupling strength as $d > \frac{\alpha_{\min}}{\lambda}$ ensures the stability of the synchronization manifold. In other words, the synchronization region, which is unbounded depends only on λ_2 . In class Γ_2 (Fig. 1c), the master stability function versus α has two zero-crossing points, α_{\min} and α_{\max} , where the region $\alpha_{\min} < \alpha < \alpha_{\max}$ is the stability region ($\Lambda < 0$). Therefore, an upper bound of the eigenvalues is also required for the stability region. In this case, the synchronization is stable if

$$\alpha_{\min} < d\lambda_2 \le d\lambda_3 \le \ldots \le d\lambda_N < \alpha_{\max} \tag{4}$$

Consequently, synchronization can be achieved for $\frac{\alpha_{\min}}{\lambda_2}$ 102 $d < \frac{\alpha_{\max}}{\lambda_N}$. By taking $R \equiv \frac{\lambda_N}{\lambda_2}$ as an eigenratio, the synchronization can occur if $R < \frac{\alpha_{\max}}{\alpha_{\min}}$. Thus, in this class, 103 104



Fig. 1: Different classes of master stability function. **a**) Class Γ_0 with no zero-crossing point, **b**) class Γ_1 with only one zero-crossing point, **c**) class Γ_2 with two zero-crossing points, and **d**) class Γ_3 with three zero-crossing.

the stability region of synchronization depends only on
the value of *R*. Barahona and Pecora [35] investigated
the stability of synchronization in small-world networks
by using the concept of the first-non-zero and maximum
eigenvalues of the Laplacian matrix.

Finally, the fourth class belongs to the master stability function with more than two zero-crossing points; as an example, class Γ_3 with three zero-crossing is illustrated in Fig. 1d. For these systems, the synchronization can be achieved if all $d\lambda_i$ reside in the $\Lambda < 0$ regions. Since this class is more complex and case-dependent, we ignore it in this study.

According to the above definitions of the master stabil-117 ity function classifications, the synchronization region is 118 only affected by λ_2 and λ_N . In fact, two networks have 119 the same synchronization region if they have the same 120 λ_2 and λ_N . Based on this concept, a reduced network 121 can have the same synchronization pattern as the original 122 network by choosing its λ_2 and λ_{\max} the same as the 123 original network. To find the connectivity matrix with de-124 fined eigenvalues, the eigen-decomposition approach can 125 be used which is explained in the next subsection. 126

Eigen-decomposition and Gram-Schmidt orthogonalization of Laplacian matrix. Consider λ_i , i = 1, 2, ..., Nand λ'_i , i = 1, 2, ..., n as the i^{th} eigenvalue of the original and reduced Laplacian matrix, respectively, and R and R'as their eigenratio as well. To have the same synchronization pattern, we must keep $\lambda_2 \cong \lambda'_2$ and also $R \cong R'$, leading to $\lambda_N \cong \lambda'_n$. To determine Laplacian matrix of the reduced network with desired eigenvalues, the eigendecomposition factorization can be utilized. According to this factorization, any positive semidefinite matrix, e.g., A, can be factorized as

$$A = QDQ^{-1} \tag{5}$$

in which D is a diagonal matrix whose diagonal elements are the eigenvalues of A, and the corresponding eigenvectors lie in the columns of Q. Therefore, by considering Das the matrix of eigenvalues of the reduced matrix $(n \times n)$ and finding an appropriate eigenvector matrix (Q), the Laplacian matrix $A_{n \times n}$ can be computed using Eq. (5). Since the matrix A is assumed symmetric, we can write,

$$A = A^{T} = (QDQ^{-1})^{T} = (Q^{-1})^{T} DQ^{T}$$
(6)

leading to $Q^{-1} = Q^T$, where T denotes the transposed 127 matrix. Therefore, Q must be an orthogonal matrix. To 128 form an orthogonal basis, the Gram-Schmidt process can 129 be used [see Appendix for more details]. Since the first 130 eigenvalue of A is zero, its corresponding eigenvector must 131 be chosen as $v_1 = [1, 1, ..., 1]_{1 \times n}$ for the Gram-Schmidt process. Selecting the other independent basis vectors is 132 133 arbitrary. Then, using the orthogonal basis vectors, $Q_{n \times n}$ 134 can be obtained. 135

In order to determine $D_{n \times n}$, n eigenvalues in the as-136 cending order are needed, where three of them are known: 137 $\lambda'_1 = 0, \ \lambda'_2 = \lambda_2, \ \text{and} \ \lambda'_n = \lambda_N.$ The rest of the needed 138 eigenvalues (n-3 eigenvalues) are found by partitioning 139 the N-3 eigenvalues of the Laplacian matrix of the orig-140 inal network. Here, we use the k-means clustering algo-141 rithm. K-means is the most popular clustering method 142 due to its simplicity (for more detail, see [36]). After ob-143 taining an orthogonal matrix Q, and a diagonal matrix 144 D, a Laplacian matrix A with the desired dimension and 145 eigenvalues can be found by using Eq. (5). It should be 146 noted that the obtained matrix is weighted. The described 147 method for obtaining the reduced connectivity matrix A148 is presented in Fig. 2. 149

Simulation results. – In this section, we apply the proposed method to reduce a high-dimensional Watts-Strogatz small-world network with N = 500 nodes and 10^5 links. It is assumed that the individual dynamics of the node obey the chaotic Sprott-B equations [37]:

$$\begin{cases} \dot{x} = yz \\ \dot{y} = x - y \\ \dot{z} = 1 - xy \end{cases}$$
(7)

The size of the reduced network is assumed as n = 100150 here. We consider different coupling functions to inves-151 tigate different synchronization patterns. For the orig-152 inal network, we have $\lambda_2 = 339.47$ and $\lambda_N = 449.80$. 153 Thus, we keep these eigenvalues and obtain the other 154 n-3 eigenvalues by classifying N-3 eigenvalues of the 155 original network. So, the matrix D is found. Next, the 156 eigen-decomposition factorization and Gram-Schmidt or-157 thogonalization are employed, and an orthogonal matrix 158 of eigenvectors is obtained (Q). Finally, a zero-row sum, 159 symmetry Laplacian matrix of size n = 100 with desired 160 eigenvalues is found using Eq. (5). The values of the 161 two most essential eigenvalues and eigenratio used in this 162 example are represented in Table 1. It can be seen that 163



Fig. 2: The schematic of the proposed method to reduce an N-dimensional network to n-dimensional one (N > n) using eigen-decomposition factorization and Gram-Schmidt orthogonalization.

Table 1: Two eigenvalues and eigenratios of the reduced network and its parent.

	Original network	Reduced network
λ_2	339.47	339.50
$\lambda_{ m max}$	449.80	449.80
R	1.32	1.32

the eigenvalues of the reduced and original networks areapproximately equal.

For more investigation, three couplings with different 166 MSF classes are considered. In Fig. 3, the master stability 167 functions versus α are plotted. Three different couplings 168 $y \to x, x \to y$, and $x \to z$ are considered. The notation, 169 e.g., $x \to z$, means that the coupling which is defined on x 170 state variables is added to z state variables. According to 171 the eigenvalues presented in Table 1, the stability regions 172 in $y \to x$ coupling are d > 0.003069 and d > 0.003081 for 173 the original and reduced networks, respectively. For $x \to y$ 174 coupling, the stability regions of the original and reduced 175 networks are 0.002957 < d < 0.0031030 and 0.002962 <176 d < 0.0031023, respectively. 177

Next, the networks are solved numerically, and the synchronization error is calculated using Eq. 8.

$$Err = \frac{1}{T(N-1)} \lim_{T \to \infty} \int_0^T \sum_{k=2}^N \sqrt{(x_1 - x_k)^2 + (y_1 - y_k)^2 + (z_1 - z_k)^2} \, \mathrm{d}t$$
(8)

The synchronization errors for both networks and each coupling scheme are illustrated in Fig. 4. The upper and lower panels represent the errors of the parent and reduced networks, respectively. It can be observed that the synchronization regions, i.e., the region of coupling strength (d) with zero error, are the same for both networks. Moreover, the synchronization errors have similar trends in the original and reduced networks.

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To better compare the synchronization behavior of both 186 networks, time series, spatiotemporal patterns, and time 187 snapshots are presented in Figs. 5-8 for synchronous and 188 asynchronous states for master stability function of class 189 Γ_1 and class Γ_2 . Figure 5 illustrates the patterns of both 190 networks for $d = 3.7 \times 10^{-3}$ which is in the synchroniza-191 tion regime under $y \rightarrow x$ coupling. Also, the results 192 for $d = 2.7 \times 10^{-3}$ in which the oscillators of networks 193 under $y \to x$ coupling are asynchronous, are shown in 194 Fig. 6. Moreover, the networks have the same behavior 195 for class Γ_2 ($x \to y$ coupling). In Fig. 7 and Fig. 8, 196 the synchronous and asynchronous behavior of both net-197 works is represented by considering $d = 3.0 \times 10^{-3}$ and 198 $d = 3.2 \times 10^{-3}$, respectively. It can be observed that the 199 networks have similar synchronous and asynchronous pat-200 terns. 201

Conclusion. – Large-scale complex networks are im-202 portant models for describing various real-world networks. 203 However, their high dimensionality often gives rise to 204 high computational costs for analysis, leading to be time-205 consuming. Hence, reducing the dimension of these net-206 works is essential. On the other hand, synchronization is a 207 significant phenomenon in complex networks. Therefore, 208 it is desired not to disturb the synchronization pattern 209 during dimension reduction. This study addressed this 210 issue by decreasing the size of the Laplacian matrix of a 211 large-scale network using the eigen-decomposition method 212 and the Gram-Schmidt orthogonalization process. The 213 original network is considered to be undirected; there-214 fore, the eigenvalues of the Laplacian matrix are real. To 215 construct a network with eigen-decomposition approach, 216 firstly, the eigenvalues of the reduced Laplacian matrix 217



Fig. 3: The master stability function versus α for Sprott-B chaotic system (Eq. (7)) under three different couplings: **a**) $y \to x$ (class Γ_1), **b**) $x \to y$ (class Γ_2), and **c**) $x \to z$ (class Γ_0). Coupled Sprott-B systems represent different synchronization patterns according to the coupling scheme.



Fig. 4: The synchronization errors of coupled Sprott-B systems for the original (upper plots) and reduced (lower plots) networks as a function of coupling strength d. The coupling is on a) class Γ_1 ($y \to x$), b) class Γ_2 ($x \to y$), and c) class Γ_0 ($x \to z$). The synchronization region and the trend of error are similar for both networks in each class.



Fig. 5: a) Time series, b) spatiotemporal pattern, and c) time snapshot at t = 4000 for $y \to x$ coupling which is class Γ_1 . The left and right panels are the results of the original network (N = 500) and the reduced one (n = 100), respectively. The coupling strength is $d = 3.7 \times 10^{-3}$, in which all oscillators lie in the synchronous manifold. The oscillations of both original and reduced networks are synchronous in this case.

must be defined. According to the master stability func-218 tion, the region of stable synchronization depends on the 219 minimum and maximum non-zero eigenvalues. Thus, we 220 kept them the same as the original network and selected 221 the other eigenvalues by classifying the original eigenval-222 ues. Then, the matrix of eigenvectors was obtained by 223 the Gram-Schmidt orthogonalization process. Finally, us-224 ing the eigenvalues and eigenvectors, a weighted reduced 225 Laplacian matrix was obtained. The method was applied 226 on a 500-node small-world network of Sprott-B systems. 227 The results were validated via synchronization error, time 228 series, spatiotemporal patterns, and snapshots of both net-229 works for different coupling functions in the synchronous 230 and asynchronous states. Our findings indicate that the 231 number of nodes of any complex network can be decreased 232 regardless of network topology and node dynamics with 233 preserving the synchronization stability region. 234

Conflict of interest. – The authors declare that they 235 have no conflict of interest. 236

APPENDIX: the Gram-Schmidt process. -237 Suppose the arbitrary set $\{\vec{v}_1, \vec{v}_2, \dots, \vec{v}_k\}$ as the ba-238 sis for a given set V, whose vectors are linearly inde-239 pendent. The Gram-Schmidt process can generate an or-240 thogonal basis for V. The vectors $\{\vec{u}_1, \vec{u}_2, \dots, \vec{u}_k\}$ are 241 said to be orthogonal if and only if the inner product of 242 any two different vectors of them is equal to zero, i.e., 243 $\langle \vec{u}_i, \vec{u}_j \rangle = 0 \ \forall \ i \neq j$. This set of new vectors can be 244 constructed as follows: 245



Fig. 6: a) Time series, b) spatiotemporal pattern, and c) time snapshot at t = 4000 for $y \to x$ coupling which is class Γ_1 . The coupling strength is $d = 2.7 \times 10^{-3}$, that leads to asynchronous oscillations in the original (left panel) and the reduced networks (right panel). This case exhibits asynchronous oscillations in both networks.

$$\begin{cases} \overrightarrow{u}_1 &= \overrightarrow{v}_1 \\ \overrightarrow{u}_2 &= \overrightarrow{v}_2 - \frac{\langle \overrightarrow{v}_2, \overrightarrow{u}_1 \rangle}{\langle \overrightarrow{u}_1, \overrightarrow{u}_1 \rangle} \overrightarrow{u}_1 \\ \overrightarrow{u}_3 &= \overrightarrow{v}_3 - \frac{\langle \overrightarrow{v}_3, \overrightarrow{u}_1 \rangle}{\langle \overrightarrow{u}_1, \overrightarrow{u}_1 \rangle} \overrightarrow{u}_1 - \frac{\langle \overrightarrow{v}_3, \overrightarrow{u}_2 \rangle}{\langle \overrightarrow{u}_2, \overrightarrow{u}_2 \rangle} \overrightarrow{u}_2 \\ \vdots \\ \overrightarrow{u}_k &= \overrightarrow{v}_k - \sum_{p=1}^{k-1} \frac{\langle \overrightarrow{v}_k, \overrightarrow{u}_p \rangle}{\langle \overrightarrow{u}_p, \overrightarrow{u}_p \rangle} \overrightarrow{u}_p \end{cases}$$

where $\langle . \rangle$ denotes the inner product.

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Fig. 7: a) Time series, b) spatiotemporal pattern, and c) time snapshot at t = 4000 for $x \to y$ coupling which is class Γ_2 for the original (left panel) and the reduced networks (right panel) at $d = 3.0 \times 10^{-3}$. Synchronized oscillations are observed in both networks.

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Fig. 8: a) Time series, b) spatiotemporal pattern, and c) the last time snapshot at t = 4000 for $x \to y$ coupling, which is class Γ_2 , with $d = 3.2 \times 10^{-3}$. Asynchronous oscillations are observed in both the left panel (the original network) and the right panel (the reduced network). It appears that in this case both networks oscillate asynchronously.

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